

HW 05 Solutions

①

4.4.1 Vibrating string of density ρ_0 , tension T_0

(a) Fixed ends: $u(0,t) = u(L,t) = 0$
 $u_{tt} = c^2 u_{xx}$, $c^2 = \frac{T_0}{\rho_0}$

Natural frequencies are $c\sqrt{\lambda}$ and

$$\lambda = \left(\frac{n\pi}{L}\right)^2 \text{ according to section 4.4.}$$

$$\Rightarrow \text{natural frequencies} = n \frac{\pi c}{L}, n=1, 2, \dots$$

(b) $u(0,t) = \frac{du(H)}{dx} = 0$

$$u = \phi(x) h(t)$$

ODE BVP
$$\begin{cases} \frac{d^2 \phi}{dx^2} = -\lambda \phi & \Rightarrow \phi = c_1 \cos(\sqrt{\lambda}x) + c_2 \sin(\sqrt{\lambda}x) \\ \phi(0) = 0 & \phi(0) = c_1 = 0 \\ \frac{d\phi(H)}{dx} = 0 & \Rightarrow \phi(x) = c_2 \sin(\sqrt{\lambda}x) \\ & \frac{d\phi}{dx} = c_2 \sqrt{\lambda} \cos(\sqrt{\lambda}x) \end{cases}$$

$$\frac{d\phi(H)}{dx} = c_2 \sqrt{\lambda} \cos(\sqrt{\lambda}H) = 0$$

$\lambda < 0$, $\lambda = 0$ - only trivial sol

$$\Rightarrow \cos(\sqrt{\lambda}H) = -\frac{\sqrt{\lambda}}{\sqrt{\lambda}} + \sqrt{\lambda}n$$

$$\Rightarrow \lambda = \frac{\pi^2 \left(n + \frac{1}{2}\right)^2}{H^2}, n=1, 2, \dots$$

$$h'' = -dc^2 h$$

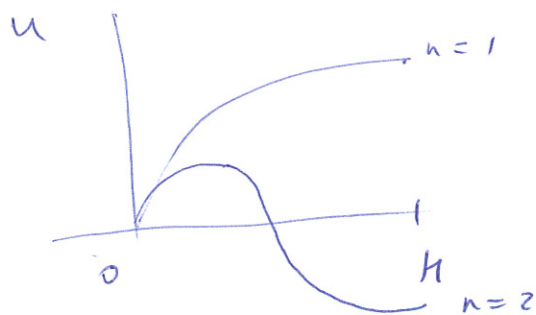
②

$$\Rightarrow h = c_1 \cos(c\sqrt{\lambda}t) + c_2 \sin(c\sqrt{\lambda}t)$$

$c\sqrt{\lambda}$ - natural frequencies

$$\Rightarrow c\sqrt{\lambda} = \frac{(n - \frac{1}{2})\pi}{H} c, \quad n = 1, 2, \dots$$

eigenfunctions $\sin\left(\frac{(n - \frac{1}{2})\pi x}{H}\right)$



(c) Odd harmonics for u !

$$u = \Phi(x) h(t)$$

$$H = \frac{L}{2}$$

$$\begin{cases} \frac{d^2 h}{dt^2} = -\lambda h \\ \frac{d^2 \Phi}{dx^2} = -\lambda \Phi \\ \Phi(0) = \Phi(L) = 0 \end{cases}$$

$$\lambda = \left(\frac{n\pi}{L}\right)^2$$

For n odd: $n = 1, 3, 5, \dots$

$$\lambda = \left(\frac{(2m-1)\pi}{L}\right)^2 = \frac{(m - \frac{1}{2})^2 \pi^2}{(\frac{L}{2})^2} = \frac{(m - \frac{1}{2})^2 \pi^2}{H^2}$$

$m = 1, 2, \dots$

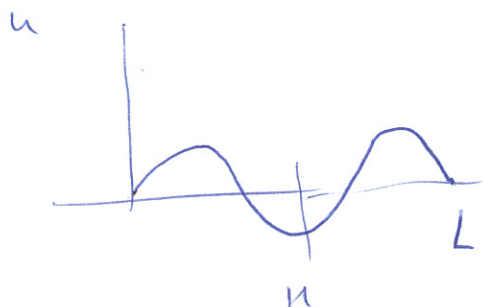
\Rightarrow frequencies of (a) = (b) for $n = 2m - 1, m = 1, 2, \dots$

Symmetry: for n odd modes we have $n - 1$,

i.e. even number of zeros

Symmetric relative to $x = L$.

$$\Rightarrow \frac{du}{dx}(L) = 0.$$



4.4.3

$$\rho_0 \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} - \beta \frac{\partial u}{\partial t}.$$

(a) β is positive hence damping

slows down motion. E.g. for $\frac{\partial u}{\partial x} = 0$

$$\Rightarrow \rho_0 \frac{d^2 u}{dt^2} = -\beta \frac{du}{dt}$$

$$\Rightarrow \frac{du}{dt} = \frac{du}{dt}(0) e^{-\frac{\beta}{\rho_0} t} \quad \text{-- decays for } \beta > 0 \text{ hence } \rho_0 > 0.$$

(b) BC $u(0, t) = u(L, t) = 0$

IC $u(x, 0) = f(x)$ and $\frac{\partial u}{\partial x}(x, 0) = g(x)$

Also assume $\beta^2 < 4\pi^2 \frac{\rho_0 T_0}{L^2}$

Separation of variables:

$$u = \phi(x) h(t)$$

$$\Rightarrow \rho_0 \phi \frac{d^2 h}{dt^2} = T_0 h \frac{d^2 \phi}{dx^2} - \beta \phi \frac{dh}{dt}$$

$$T_0 \phi h$$

$$\Rightarrow \frac{\rho_0}{T_0} \frac{d^2 h}{dt^2} + \frac{\beta}{T_0} \frac{dh}{dt} = \frac{d^2 \phi}{dx^2} = -\lambda$$

ODE BVP: $\begin{cases} \frac{d^2 \phi}{dx^2} = -\lambda \phi \\ \phi(0) = \phi(L) = 0 \end{cases} \Rightarrow \lambda = \left(\frac{n\pi}{L}\right)^2$
 $\Rightarrow \phi = \sin\left(\frac{n\pi x}{L}\right),$
 $n = 1, 2, \dots$

Equation for h:

$\rho_0 h'' + \beta h' + T_0 \lambda h = 0$ — homogeneous
 linear ODE with constant coefficients
 \Rightarrow solve in the form $h = e^{rt}$

\Rightarrow characteristic eq:

$$\rho_0 r^2 + \beta r + \left(\frac{n\pi}{L}\right)^2 T_0 = 0$$

$$\Rightarrow r_{1,2} = \frac{-\beta \pm \sqrt{\beta^2 - 4\rho_0 T_0 \left(\frac{h\pi}{L}\right)^2}}{2\rho_0} \quad (5)$$

Since $\beta^2 < 4\rho_0 T_0 \left(\frac{h\pi}{L}\right)^2 \Rightarrow$ discriminant $< 0 \forall n$

\Rightarrow complex eigenvalue:

$$r_{1,2} = -\frac{\beta}{2\rho_0} \pm i \sqrt{\frac{T_0}{\rho_0} \left(\frac{h\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}}$$

General sol:

$$h = c_1 e^{-\frac{\beta}{2\rho_0} t} \cos \left[\sqrt{\frac{T_0}{\rho_0} \left(\frac{h\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}} t \right]$$

$$+ c_2 e^{-\frac{\beta}{2\rho_0} t} \sin \left[\sqrt{\frac{T_0}{\rho_0} \left(\frac{h\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}} t \right]$$

Here we also see explanation for (a) that for $\beta < 0$ we would have exp growing solution which is unphysical for damping.

Superposition:

$$u = e^{-\frac{\beta}{2\rho_0} t} \sum_{n=1}^{\infty} \left[a_n \cos(\omega_n t) + b_n \sin(\omega_n t) \right] \sin \frac{n\pi x}{L}$$

where $\omega_n = \sqrt{\frac{T_0}{\rho_0} \left(\frac{h\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}}$

$$IC : u(x, 0) = f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)$$

$$\Rightarrow a_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$\frac{du}{dt}(x, 0) = g(x) = -\frac{\beta}{2\rho_0} e^{-\frac{\beta t}{2\rho_0}} \sum_{n=1}^{\infty} (a_n \cos(\omega_n t))$$

$$+ b_n \sin(\omega_n t) \Big|_{t=0} + e^{-\frac{\beta t}{2\rho_0}} \sum_{n=1}^{\infty} \omega_n (a_n \sin(\omega_n t))$$

$$+ b_n \cos(\omega_n t) \Big] \sin\left(\frac{n\pi x}{L}\right) \Big|_{t=0}$$

$$= -\frac{\beta}{2\rho_0} \underbrace{\sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)}_{f(x)} + \sum_{n=1}^{\infty} b_n \omega_n \sin\left(\frac{n\pi x}{L}\right)$$

$$\Rightarrow b_n \omega_n = \frac{\beta a_n}{2\rho_0} + \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

already given above

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5.3.3 $H(x) \left| \frac{d^2 \phi}{dx^2} + \alpha(x) \frac{d\phi}{dx} + [\lambda \beta(x) + \gamma(x)] \phi = 0 \quad (*) \right.$

We need to find $H(x)$ such that the resulting eq is in SL form:

$$\frac{d}{dx} \left[p(x) \frac{d\phi}{dx} \right] + [\lambda \delta(x) + \epsilon(x)] \phi = 0 \quad (**)$$

$$H(x) \frac{d^2 \phi}{dx^2} + H(x) \alpha(x) \frac{d\phi}{dx} = \frac{d}{dx} \left[p(x) \frac{d\phi}{dx} \right]$$

$$= p' \frac{d\phi}{dx} + p \frac{d^2 \phi}{dx^2}$$

\Rightarrow terms with $\frac{d^2 \phi}{dx^2}$ $\left. \begin{array}{l} p'(x) = H(x) \alpha(x) \\ p(x) = H(x) \end{array} \right\} \Rightarrow H' = H \alpha$

$$\frac{dH}{H} = \alpha(x) dx$$

$$\Rightarrow H(x) = c_1 e^{\int \alpha(x) dx}$$

Set e.g. $c_1 = 1 \Rightarrow H = e^{\int \alpha(x) dx}$

compare (*) and (**): $\Rightarrow p(x) = H(x), \quad \epsilon(x) = \gamma H, \quad \delta = \beta H.$

5.3.5

(2)

$$\phi'' + \lambda \phi = 0, \quad \frac{d\phi(0)}{dx} = \frac{d\phi(L)}{dx} = 0$$

(8)

$$\Rightarrow \phi = c_1 \cos(\sqrt{\lambda} x) + c_2 \sin(\sqrt{\lambda} x)$$

$$\frac{d\phi}{dx} = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda} x) + c_2 \sqrt{\lambda} \cos(\sqrt{\lambda} x)$$

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$$\frac{d\phi(0)}{dx} = 0 = c_2 \sqrt{\lambda} \Rightarrow c_2 = 0$$

$$\frac{d\phi(L)}{dx} = \frac{d\phi}{dx}(L) = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda} L) = 0$$

$\lambda = 0$ $\phi = \text{const}$

$$\Rightarrow \lambda = \left(\frac{n\pi}{L}\right)^2, \quad \phi_n = \cos\left(\frac{n\pi}{L} x\right), \quad n=0, 1, \dots$$

(a) ∞ of eigenvalues.

(b) $\cos\left(\frac{n\pi}{L} x\right)$ has $n-1$ zeros at $0 < x < L$

(c) Fourier cosine series has use of orthogonality of cosines. Completeness - from thm. on Fourier Series.

not assigned

(d) Rayleigh quotient

$$\lambda = \frac{\phi \frac{d\phi}{dx} \Big|_0^L + \int_0^L \left(\frac{d\phi}{dx}\right)^2 dx}{\int_0^L \phi^2 dx} \geq 0 \Rightarrow \lambda < 0 \text{ is not eigenval}$$

$\lambda = 0$ is eigenvalue for $\phi = \text{const}$.

$$x^2 \frac{d^2 \phi}{dx^2} + x \frac{d\phi}{dx} + \lambda \phi = 0$$

$$\phi(1) = 0, \quad \phi(b) = 0$$

Since eq. is equidimensional $\Rightarrow \phi = x^r$

$$\Rightarrow r(r-1) + r + \lambda = 0 \Rightarrow r^2 = -\lambda$$

$$\underline{\lambda > 0} \quad r = \pm i\sqrt{\lambda} \Rightarrow x^r = x^{\pm i\sqrt{\lambda}} = e^{\pm i\sqrt{\lambda} \ln x}$$

\Rightarrow general sol $\phi = C_1 \cos(\sqrt{\lambda} \ln x) + C_2 \sin(\sqrt{\lambda} \ln x)$

$$\underline{BC} \quad \phi(1) = 0 = C_1$$

$$\phi(b) = C_2 \sin(\sqrt{\lambda} \ln b) = 0 \Rightarrow \sqrt{\lambda} \ln b = n\pi, \quad n=1, 2, \dots$$

$$\Rightarrow \lambda = \left(\frac{n\pi}{\ln b} \right)^2$$

Eigen function $\sin\left(\frac{n\pi}{\ln b} \ln x\right)$

$$\underline{\lambda = 0} \Rightarrow x^2 \frac{d^2 \phi}{dx^2} + x \frac{d\phi}{dx} = x \frac{d}{dx} \left(x \frac{d\phi}{dx} \right) = 0$$

$$\Rightarrow \frac{d}{dx} \left(x \frac{d\phi}{dx} \right) = 0$$

$$x \frac{d\Phi}{dx} = c_1$$

$$d\Phi = c_1 \frac{dx}{x}$$

$$\Phi = c_1 \ln x + c_2$$

BC

$$\Phi(1) = 0 = c_2$$

$$\Phi(b) = c_1 \ln b = 0 \Rightarrow c_1 = 0 \Rightarrow \text{not a non-trivial solution}$$

$\lambda < 0$ ~~not correct to analyze~~. $r = e^{+\sqrt{\lambda} \ln x}$
general sol

$$\Phi = c_1 e^{\sqrt{\lambda} \ln x} + c_2 e^{-\sqrt{\lambda} \ln x}$$

$$\Phi(1) = c_1 + c_2 = 0 \Rightarrow c_2 = -c_1$$

$$\Phi(b) = c_1 (e^{\sqrt{\lambda} \ln b} - e^{-\sqrt{\lambda} \ln b}) = 0$$

$\Rightarrow c_1 = 0$ - only trivial sol

So $\lambda_n = \left(\frac{n\pi}{\ln b}\right)^2 \rightarrow \infty \Rightarrow n \rightarrow \infty$

Smallest λ : $\lambda_1 = \left(\frac{\pi}{\ln b}\right)^2$