

HW 05 Solutions

①

4.4.5

Vibrating string of density ρ_0 , tension T_0

(a) Fixed ends: $u(0,t) = u(L,t) = 0$

$$u_{tt} = c^2 u_{xx}, \quad c^2 = \frac{T_0}{\rho_0}$$

Natural frequencies are $c\sqrt{\lambda}$ and

$$\lambda = \left(\frac{n\pi}{L}\right)^2 \text{ according to section 4.4.}$$

$$\Rightarrow \text{natural frequencies} = n \frac{\pi c}{L}, n=1, 2, \dots$$

(b) $u(0,t) = \frac{du(H)}{dx} = 0$

$$u = \phi(x) h(t)$$

ODE BVP
$$\begin{cases} \frac{d^2 \phi}{dx^2} = -\lambda \phi & \Rightarrow \phi = c_1 \cos(\sqrt{\lambda}x) + c_2 \sin(\sqrt{\lambda}x) \\ \phi(0) = 0 & \Rightarrow c_1 = 0 \\ \frac{d\phi(H)}{dx} = 0 & \Rightarrow \phi(x) = c_2 \sin(\sqrt{\lambda}x) \end{cases}$$

$$\frac{d\phi}{dx} = c_2 \sqrt{\lambda} \cos(\sqrt{\lambda}x)$$

$$\frac{d\phi(H)}{dx} = c_2 \sqrt{\lambda} \cos(\sqrt{\lambda}H) = 0$$

$\lambda < 0$, $\lambda = 0$ - only trivial sol

$$\Rightarrow \cancel{H\sqrt{\lambda}} = -\frac{\pi}{2} + \pi n$$

$$\Rightarrow \lambda = \frac{\pi^2 (n + \frac{1}{2})^2}{H^2}, n=1, 2, \dots$$

$$h'' = -dc^2 h$$

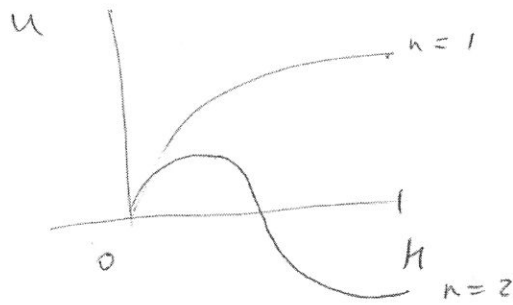
(2)

$$\Rightarrow h = c_1 \cos(c\sqrt{\lambda}t) + c_2 \sin(c\sqrt{\lambda}t)$$

$c\sqrt{\lambda}$ - natural frequencies

$$\Rightarrow c\sqrt{\lambda} = \frac{(n - \frac{1}{2})\pi}{H} c, \quad n = 1, 2, \dots$$

eigenfunctions $\sin\left(\frac{(n - \frac{1}{2})\pi x}{H}\right)$



(c) Odd harmonics for u :

$$u = \Phi(x) h(t)$$

$$H = \frac{L}{2}$$

$$\begin{cases} \frac{d^2 h}{dt^2} = -\lambda h \\ \frac{d^2 \Phi}{dx^2} = -\lambda \Phi \\ \Phi(0) = \Phi(L) = 0 \end{cases}$$

$$\lambda = \left(\frac{n\pi}{L}\right)^2$$

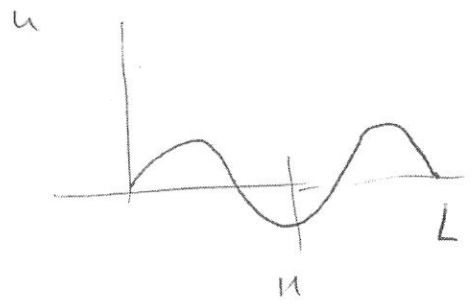
For odd n : $n = 1, 3, 5, \dots$

$$\lambda = \left(\frac{(2m-1)\pi}{L}\right)^2 = \frac{(m - \frac{1}{2})^2 \pi^2}{(L/2)^2} = \frac{(m - \frac{1}{2})^2 \pi^2}{H^2}$$

$m = 1, 2, \dots$

⇒ frequencies of (a) = (b) for $n = 2m - 1, m = 1, 2, \dots$

Symmetry: for n odd modes we have $n - 1$,
i.e. even number of zeros
Symmetric relative to $x = L$.



$$\Rightarrow \frac{du(L)}{dx} = 0.$$

4.4.3

$$\rho_0 \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} - \beta \frac{\partial u}{\partial t}$$

(a) β is positive hence damping
slows down motion. E.g. for $\frac{\partial u}{\partial x} = 0$

$$\Rightarrow \rho_0 \frac{d^2 u}{dt^2} = -\beta \frac{du}{dt}$$

$$\Rightarrow \frac{du}{dt} = \frac{du}{dt}(0) e^{-\frac{\beta}{\rho_0} t} \quad \text{— decays for } \beta > 0 \text{ hence } \rho_0 > 0.$$

(b) BC $u(0, t) = u(L, t) = 0$

IC $u(x, 0) = f(x)$ and $\frac{\partial u}{\partial x}(x, 0) = g(x)$

Also assume $\beta^2 < 4\pi^2 \rho_0 \frac{T_0}{L^2}$

Separation of variables:

(7)

$$u = \Phi(x) h(t)$$

$$\Rightarrow \rho_0 \Phi \frac{d^2 h}{dt^2} = T_0 h \frac{d^2 \Phi}{dx^2} - \beta \Phi \frac{dh}{dt}$$

$$T_0 \Phi h$$

$$\Rightarrow \frac{\rho_0}{T_0} \frac{d^2 h}{dt^2} + \frac{\beta}{T_0} \frac{dh}{dt} = \frac{d^2 \Phi}{dx^2} = -\lambda$$

ODE BVP:
$$\begin{cases} \frac{d^2 \Phi}{dx^2} = -\lambda \Phi \\ \Phi(0) = \Phi(L) = 0 \end{cases} \Rightarrow \lambda = \left(\frac{n\pi}{L}\right)^2$$

$$\Rightarrow \Phi = \sin\left(\frac{n\pi x}{L}\right),$$

$$n = 1, 2, \dots$$

Equation for h :

$$\rho_0 h'' + \beta h' + T_0 \lambda h = 0 \quad - \text{homogeneous}$$

linear ODE with constant coefficients

\Rightarrow solve in the form $h = e^{rt}$

\Rightarrow characteristic eq:

$$\rho_0 r^2 + \beta r + \left(\frac{n\pi}{L}\right)^2 T_0 = 0$$

$$\Rightarrow r_{1,2} = \frac{-\beta \pm \sqrt{\beta^2 - 4\rho_0 T_0 \left(\frac{n\pi}{L}\right)^2}}{2\rho_0} \quad (5)$$

Since $\beta^2 < 4\rho_0 T_0 \left(\frac{n\pi}{L}\right)^2 \Rightarrow$ discriminant $< 0 \forall n$

\Rightarrow complex eigenvalue:

$$r_{1,2} = -\frac{\beta}{2\rho_0} \pm i \sqrt{\frac{T_0}{\rho_0} \left(\frac{n\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}}$$

General sol:

$$h = c_1 e^{-\frac{\beta}{2\rho_0} t} \cos \left[\sqrt{\frac{T_0}{\rho_0} \left(\frac{n\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}} t \right]$$

$$+ c_2 e^{-\frac{\beta}{2\rho_0} t} \sin \left[\sqrt{\frac{T_0}{\rho_0} \left(\frac{n\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}} t \right]$$

Here we also see explanation for (a) that for $\beta < 0$ we would have exp growing solution which is unphysical for damping.

Superposition:

$$u = e^{-\frac{\beta}{2\rho_0} t} \sum_{n=1}^{\infty} \left[a_n \cos(\omega_n t) + b_n \sin(\omega_n t) \right] \sin \frac{n\pi x}{L},$$

$$\text{where } \omega_n = \sqrt{\frac{T_0}{\rho_0} \left(\frac{n\pi}{L}\right)^2 - \frac{\beta^2}{4\rho_0^2}}$$

⑥

$$IC : u(x, 0) = f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)$$

$$\Rightarrow a_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$\frac{du}{dt}(x, 0) = g(x) = -\frac{\beta}{2\rho_0} e^{-\frac{\beta t}{2\rho_0}} \sum_{n=1}^{\infty} [a_n \cos(\omega_n t)]$$

$$+ b_n \sin(\omega_n t) \Big|_{t=0} + e^{-\frac{\beta t}{2\rho_0}} \sum_{n=1}^{\infty} \omega_n (a_n \sin(\omega_n t) + b_n \cos(\omega_n t)) \sin\left(\frac{n\pi x}{L}\right) \Big|_{t=0}$$

$$= -\frac{\beta}{2\rho_0} \underbrace{\sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)}_{f(x)} + \sum_{n=1}^{\infty} b_n \omega_n \sin\left(\frac{n\pi x}{L}\right)$$

$$\Rightarrow b_n \omega_n = \frac{\beta a_n}{2\rho_0} + \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

already given above

⑦

$$\textcircled{5.3.3} \quad H(x) \left| \frac{d^2 \phi}{dx^2} + z(x) \frac{d\phi}{dx} + [\lambda \beta(x) + \gamma(x)] \phi = 0 \right.$$

We need to find $H(x)$ such that the resulting eq is in SL form:

$$\frac{d}{dx} \left[p(x) \frac{d\phi}{dx} \right] + [\lambda \delta(x) + q(x)] \phi = 0$$

$$H(x) \frac{d^2 \phi}{dx^2} + H(x) z(x) \frac{d\phi}{dx} = \frac{d}{dx} \left[p(x) \frac{d\phi}{dx} \right]$$

$$= p' \frac{d\phi}{dx} + p \frac{d^2 \phi}{dx^2}$$

$$\Rightarrow \left. \begin{array}{l} p'(x) = H(x) z(x) \\ p(x) = H(x) \end{array} \right\} \Rightarrow \left. \begin{array}{l} H' = H z \\ \frac{dH}{H} = z(x) dx \end{array} \right.$$

$$\Rightarrow H(x) = c_1 e^{\int z(x) dx}$$

$$\text{Set e.g. } c_1 = 1 \Rightarrow \underline{H = e^{\int z(x) dx}}$$

$$\Rightarrow p(x) = H(x), \quad z(x) = \gamma H, \quad \delta = \beta H.$$

5.3.5

(a)

$$\Phi'' + \lambda \Phi = 0, \quad \frac{d\Phi(0)}{dx} = \frac{d\Phi(L)}{dx} = 0$$

(8)

$$\Rightarrow \Phi = c_1 \cos(\sqrt{\lambda} x) + c_2 \sin(\sqrt{\lambda} x)$$

$$\frac{d\Phi}{dx} = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda} x) + c_2 \sqrt{\lambda} \cos(\sqrt{\lambda} x)$$

$\lambda > 0$

$$\frac{d\Phi(0)}{dx} = 0 = c_2 \sqrt{\lambda} \Rightarrow c_2 = 0$$

$$\frac{d\Phi(L)}{dx} = \frac{d\Phi}{dx}(L) = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda} L) = 0$$

$\lambda = 0$ $\Phi = \text{const.}$

$$\Rightarrow \lambda = \left(\frac{n\pi}{L}\right)^2, \quad \Phi_n = \cos\left(\frac{n\pi}{L} x\right), \quad n=0, 1, \dots$$

(a) ∞ of eigenvalues.

(b) $\cos\left(\frac{n\pi}{L} x\right)$ has $n-1$ zeros at $0 < x < L$

(c) Fourier cosine series that use orthonormality of cosines. Completeness - from thm. on Fourier Series.

not assigned

(d) Rayleigh quotient

$$\lambda = \frac{\left. \Phi \frac{d\Phi}{dx} \right|_0^L + \int_0^L \left(\frac{d\Phi}{dx}\right)^2 dx}{\int_0^L \Phi^2 dx} \geq 0 \Rightarrow \lambda < 0 \text{ is not eigenvalue}$$

$\lambda = 0$ is eigenvalue for $\Phi = \text{const.}$

5.3.9.c

9

$$x^2 \frac{d^2 \phi}{dx^2} + x \frac{d\phi}{dx} + \lambda \phi = 0$$

$$\phi(1) = 0, \quad \phi(b) = 0$$

Since eq. is equidimensional $\Rightarrow \phi = x^r$

$$\Rightarrow r(r-1) + r + \lambda = 0 \Rightarrow r^2 = -\lambda$$

$\lambda > 0$ $r = \pm i\sqrt{\lambda} \Rightarrow x^r = x^{\pm i\sqrt{\lambda}} = e^{\pm i\sqrt{\lambda} \ln x}$

\Rightarrow general sol $\phi = c_1 \cos(\sqrt{\lambda} \ln x) + c_2 \sin(\sqrt{\lambda} \ln x)$

BC $\phi(1) = 0 = c_1$

$$\phi(b) = c_2 \sin(\sqrt{\lambda} \ln b) = 0 \Rightarrow \sqrt{\lambda} \ln b = n\pi, \quad n=1, 2, \dots$$

$$\Rightarrow \lambda = \left(\frac{n\pi}{\ln b} \right)^2$$

Eigen function $\sin\left(\frac{n\pi}{\ln b} \ln x\right)$

$\lambda = 0$ $\Rightarrow x^2 \frac{d^2 \phi}{dx^2} + x \frac{d\phi}{dx} = x \frac{d}{dx} \left(x \frac{d\phi}{dx} \right) = 0$

$$\Rightarrow \frac{d}{dx} \left(x \frac{d\phi}{dx} \right) = 0$$

$$x \frac{d\Phi}{dx} = c_1$$

$$d\Phi = c_1 \frac{dx}{x}$$

$$\Phi = c_1 \ln x + c_2$$

BC

$$\Phi(1) = 0 = c_2$$

$$\Phi(b) = c_1 \ln b = 0 \Rightarrow c_1 = 0 \Rightarrow \text{not eigenvalue}$$

$\lambda < 0$ ~~not need to analyze~~. $r = e^{\pm \sqrt{\lambda} \ln x}$
general sol

$$\Phi = c_1 e^{\sqrt{\lambda} \ln x} + c_2 e^{-\sqrt{\lambda} \ln x}$$

$$\Phi(1) = c_1 + c_2 = 0 \Rightarrow c_2 = -c_1$$

$$\Phi(b) = c_1 (e^{\sqrt{\lambda} \ln b} - e^{-\sqrt{\lambda} \ln b}) = 0$$

$\Rightarrow c_1 = 0$ - only trivial sol

So $\lambda_n = \left(\frac{n\pi}{\ln b}\right)^2 \rightarrow \infty \Rightarrow n \rightarrow \infty$

Smallest λ : $\lambda_1 = \left(\frac{\pi}{\ln b}\right)^2$

math 312

HW 6 Solutions

①

5.5.1

$$p \left(u \frac{dv}{dx} - v \frac{du}{dx} \right) \Big|_a^b = 0 \Rightarrow \text{self-adjoint}$$

(c)

$$\frac{d\phi}{dx}(0) - h\phi(0) = 0, \quad \frac{d\phi}{dx}(L) = 0$$

\Rightarrow both for u and v ($a=0, b=L$):

$$\frac{du(0)}{dx} - hu(0) = 0 = \frac{dv(L)}{dx}$$

$$\frac{dv(0)}{dx} - hv(0) = 0 = \frac{du(L)}{dx}$$

$$\Rightarrow u \left(\frac{dv}{dx} - v \frac{du}{dx} \right) \Big|_0^L = \underbrace{u(L) \frac{dv(L)}{dx}}_0 - \underbrace{v(0) \frac{du(0)}{dx}}_0$$

$$- \underbrace{u(0) \frac{dv(0)}{dx}}_{hv(0)} + \underbrace{v(0) \frac{du(0)}{dx}}_{hu(0)} = -hu(0)v(0) + hv(0)u(0) = 0$$

$$(2) \quad \phi(L) + \alpha \phi(0) + \beta \frac{d\phi}{dx}(0) = 0$$

$$\frac{d\phi}{dx}(L) + \gamma \phi(0) + \delta \frac{d\phi}{dx}(0) = 0$$

Solving for $u(L), v(L); \frac{du(L)}{dx}, \frac{dv(L)}{dx}$

u: $u(L) = -\alpha u(0) - \beta \frac{du(0)}{dx}$

$\frac{du(L)}{dx} = -\gamma u(0) - \delta \frac{du(0)}{dx}$

v: $v(L) = -\alpha v(0) - \beta \frac{dv(0)}{dx}$

$\frac{dv(L)}{dx} = -\gamma v(0) - \delta \frac{dv(0)}{dx}$

($a=0, b=L$)

$\Rightarrow \left(u \frac{dv}{dx} - v \frac{du}{dx} \right) \Big|_0^L$

$= u(L) \frac{dv(L)}{dx} - u(0) \frac{dv(0)}{dx} - v(L) \frac{du(L)}{dx} + v(0) \frac{du(0)}{dx}$

$= \left\{ \text{using } (*) \right\}$

$= \left(-\alpha u(0) - \beta \frac{du(0)}{dx} \right) \left(-\gamma v(0) - \delta \frac{dv(0)}{dx} \right) - u(0) \frac{dv(0)}{dx}$

$- \left(-\alpha v(0) - \beta \frac{dv(0)}{dx} \right) \left(-\gamma u(0) - \delta \frac{du(0)}{dx} \right) + v(0) \frac{du(0)}{dx}$

(*)

(*)

(3)

$$\begin{aligned}
&= \underbrace{\alpha \gamma u(0) v(0)} + \underbrace{\beta \gamma \frac{d u(0)}{dx} v(0)} + \underbrace{\alpha \delta u(0) \frac{d v(0)}{dx}} + \underbrace{\beta \delta \frac{d u(0)}{dx} \frac{d v(0)}{dx}} \\
&- \underbrace{u(0) \frac{d v(0)}{dx}} \\
&- \underbrace{\alpha \gamma v(0) u(0)} - \underbrace{\beta \gamma \frac{d v(0)}{dx} u(0)} - \underbrace{\alpha \delta v(0) \frac{d u(0)}{dx}} - \underbrace{\beta \delta \frac{d v(0)}{dx} \frac{d u(0)}{dx}} \\
&+ \underbrace{v(0) \frac{d u(0)}{dx}}
\end{aligned}$$

$$= \frac{d u(0)}{dx} v(0) [\beta \gamma - \alpha \delta + 1] + u(0) \frac{d v(0)}{dx} [\alpha \delta - 1 - \beta \gamma] = 0$$

$$\Rightarrow \boxed{\alpha \delta - \beta \gamma - 1 = 0}$$

(5.5.5)

$$L = \frac{d^2}{dx^2} + 6 \frac{d}{dx} + 9$$

(a)

$$\begin{aligned}
L(e^{rx}) &= r^2 e^{rx} + r \cdot 6 e^{rx} + 9 e^{rx} \\
&= (r+3)^2 e^{rx}
\end{aligned}$$

(b) $L(y) = 0$ using part (a): $r = -3$ double root

$$\Rightarrow y = (c_1 + c_2 x) e^{-3x}$$

5.8.3

(a)

$$\frac{d^2 \phi}{dx^2} + \lambda \phi = 0, \quad \frac{d\phi}{dx}(0) = 0$$

(2)

$$\frac{d\phi}{dx}(L) + h\phi(L) = 0$$

$h > 0$

$$\lambda = \frac{-p\phi \frac{d\phi}{dx} \Big|_0^L + \int_0^L \left[p \left(\frac{d\phi}{dx} \right)^2 - 2\phi^2 \right] dx}{\int_0^L \phi^2 \sigma dx}$$

$$= \left. \begin{array}{l} p=1, \sigma=1, q=0 \\ \frac{d\phi}{dx}(L) = -h\phi(L) \end{array} \right\}$$

$$= \frac{h\phi(L)^2 + \int_0^L \left(\frac{d\phi}{dx} \right)^2 dx}{\int_0^L \phi^2 dx} \geq 0 \quad \text{since } h > 0$$

Also if to have $= 0$ require $\phi(L) = 0, \frac{d\phi}{dx} = 0$
 $\Rightarrow \phi(x) \equiv 0$ - only

trivial sol $\Rightarrow \lambda > 0$

(b) By (a) $\lambda > 0$

$$\Rightarrow \phi = c_1 \cos(\sqrt{\lambda}x) + c_2 \sin(\sqrt{\lambda}x)$$

$$\frac{d\Phi}{dx} = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda}x) + c_2 \sqrt{\lambda} \cos(\sqrt{\lambda}x) \quad (3)$$

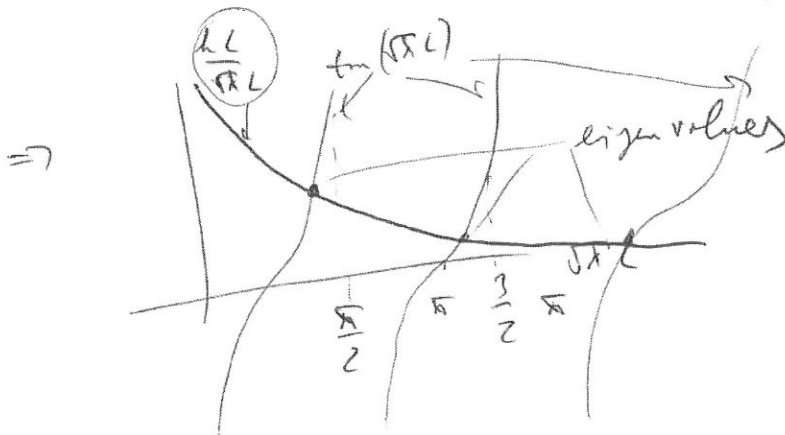
$$\frac{d\Phi}{dx}(0) = 0 = c_2 \sqrt{\lambda} \Rightarrow c_2 = 0$$

$$\Phi(x) = c_1 \cos(\sqrt{\lambda}x)$$

BC at $x=L$: $\frac{d\Phi}{dx}(L) + h\Phi(L) = -c_1 \sqrt{\lambda} \sin(\sqrt{\lambda}L)$

$$+ h c_1 \cos(\sqrt{\lambda}L) = 0$$

$$\Rightarrow \tan(\sqrt{\lambda}L) = \frac{h}{\sqrt{\lambda}} = \frac{hL}{\sqrt{\lambda}L}$$



$\frac{hL}{\sqrt{\lambda}L} = \text{hyperbola}$
 $\propto \frac{1}{\sqrt{\lambda}L}$

$$\propto \sqrt{\lambda}_1 L < \frac{\pi}{2}$$

$$\pi < \sqrt{\lambda}_2 L < \frac{3}{2}\pi$$

$$2\pi < \sqrt{\lambda}_3 L < \frac{5}{2}\pi$$

$$\Rightarrow (n-1)\pi < \sqrt{\lambda}_n L < (n-\frac{1}{2})\pi$$

As $n \rightarrow \infty$

$$\Rightarrow \lambda_n \rightarrow \left(\frac{\pi}{2}(n-1)\right)^2$$

Approach lower bound

$n = 1, 2, \dots$

5.8.7

7

$$\frac{d^2 \phi}{dx^2} + \lambda \phi = 0$$

$$\phi(0) = 0, \quad \phi(\pi) = 2 \frac{d\phi}{dx}(0)$$

$$(a) \int_0^{\pi} \left(u \frac{d^2 v}{dx^2} - v \frac{d^2 u}{dx^2} \right) dx$$

$$= \left. u \frac{dv}{dx} - v \frac{du}{dx} \right|_0^{\pi} - \int_0^{\pi} \left(\frac{du}{dx} \frac{dv}{dx} - \frac{dv}{dx} \frac{du}{dx} \right) dx$$

$$= u(\pi) \frac{dv(\pi)}{dx} - v(\pi) \frac{du(\pi)}{dx} = 2 \frac{du(0)}{dx} \cdot \frac{dv(\pi)}{dx} - 2 \frac{dv(0)}{dx} \frac{du(\pi)}{dx}$$

generally $\neq 0$

(for general values of $\frac{du(0)}{dx}$,

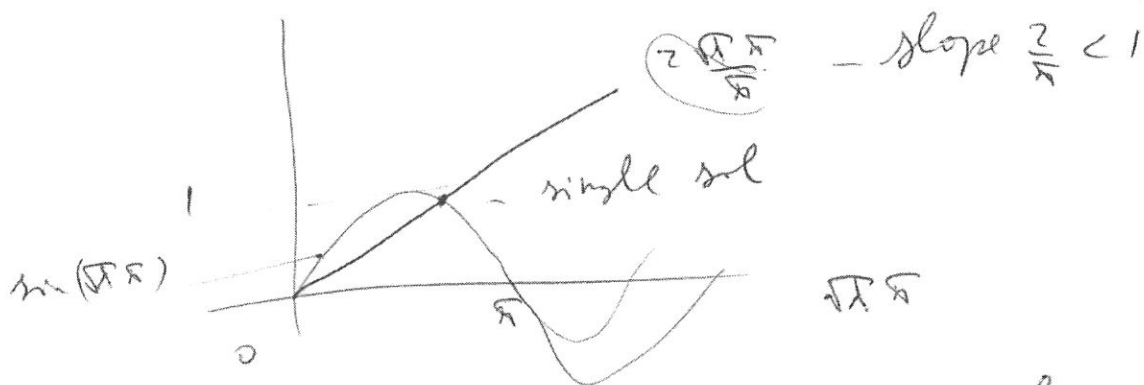
$\frac{dv(0)}{dx}$, $\frac{dv(\pi)}{dx}$, $\frac{du(\pi)}{dx}$.)

(b) If $\lambda > 0 \Rightarrow \phi(0) = 0$ results in $\phi = \sin(\sqrt{\lambda} x)$

$$\phi(\pi) - 2 \frac{d\phi}{dx} = \sin(\sqrt{\lambda} \pi) - 2\sqrt{\lambda} \cos(\sqrt{\lambda} \pi) \cdot 0 = \sin(\sqrt{\lambda} \pi) - 2\sqrt{\lambda} \cos(\sqrt{\lambda} \pi) = 0$$

Graphical sol :

(5)



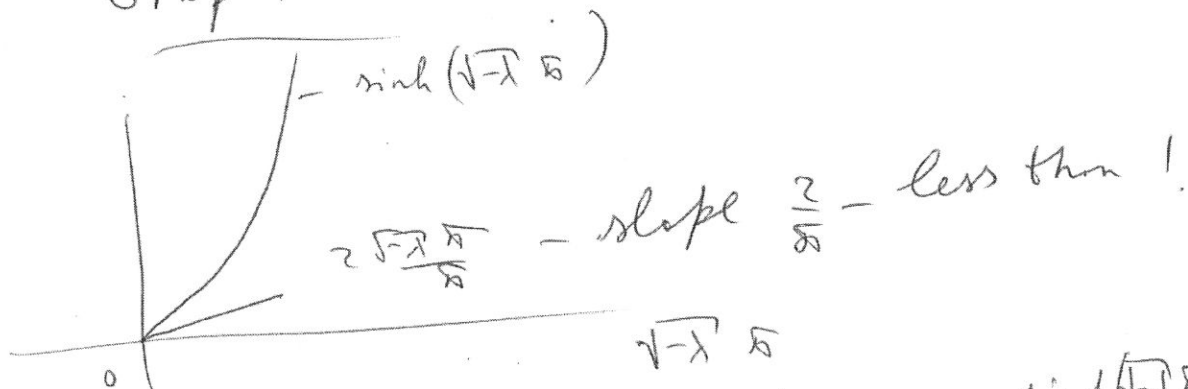
\Rightarrow 1 positive eigenvalue < 1 .

(c) If $\lambda < 0 \Rightarrow \phi(0) = 0$.

gives $\phi = \sinh(\sqrt{-\lambda}x)$

$$\phi(\pi) - 2 \frac{\phi(0)}{dx} = \sinh(\sqrt{-\lambda}\pi) - 2\sqrt{-\lambda} = 0$$

Graphical sol .



\Rightarrow no solution of $\sinh(\sqrt{-\lambda}\pi) = 2\sqrt{-\lambda}$ for $\lambda < 0$.

(6)

$$(d) \lambda = 0 \Rightarrow \phi = c_1 x + c_2$$

$$\phi(0) = c_2 = 0$$

$$\Rightarrow \phi(x) = c_1 x$$

$$\phi(\pi) - 2 \frac{d\phi(0)}{dx} = c_1 \pi - 2c_1 = 0$$

$\Rightarrow c_1 = 0$ - only trivial sol

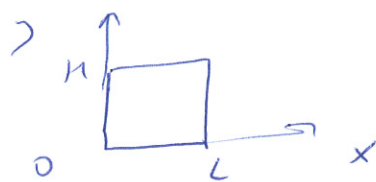
$\Rightarrow \lambda = 0$ is not eigenvalue.

(e) Since BC. are not in the form
eg 5.3, ~~2nd~~ book, then this is not

a regular Sturm-Liouville
eigenvalue problem \Rightarrow can have
complex eigenvalues!

math 312 HW solutions HW 8 ①

7.3.1



②

$$u_t = \kappa(u_{xx} + u_{yy})$$

IC: $u(x, y, 0) = d(x, y)$

$$u = \Phi(x, y) h(t) = f(x)g(y)h(t)$$

$$\frac{dh}{dt} = -\lambda h$$

$$\nabla^2 \Phi + \lambda \Phi = 0 \Rightarrow \frac{d^2 f}{dx^2} = -\mu f$$

$$\frac{d^2 g}{dy^2} + (\lambda - \mu)g = 0$$

③ $u(0, y, t) = u(L, y, t) = u(x, 0, t) = u(x, h, t) = 0$

$$\Rightarrow \begin{cases} \frac{d^2 f}{dx^2} = -\mu f \\ f(0) = f(L) = 0 \end{cases} \Rightarrow \mu = \left(\frac{h\pi}{L}\right)^2, \quad h=1, 2, \dots$$

Eigen functions:
 $f_n = \sin\left(\frac{h\pi x}{L}\right)$

(2)

$$\left\{ \begin{aligned} \frac{d^2 g}{dy^2} &= -(\lambda - \mu) g \\ g(0) &= g(H) = 0 \end{aligned} \right.$$

$$\Rightarrow \lambda - \mu = \left(\frac{n\pi}{H} \right)^2$$

Eigenfunction:

$$g_{nm} = \sin \frac{n\pi y}{H}, \quad m=1, 2, \dots$$

$$\Rightarrow \lambda_{nm} = \left(\frac{n\pi}{L} \right)^2 + \left(\frac{m\pi}{H} \right)^2$$

$$\varphi = \sin \frac{n\pi x}{L} \sin \frac{m\pi y}{H}, \quad n, m=1, 2, \dots$$

Superposition:

$$u = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} a_{nm} \sin \frac{n\pi x}{L} \sin \frac{m\pi y}{H} e^{-\lambda_{nm} t}$$

I.C. $f = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} a_{nm} \sin \frac{n\pi x}{L} \sin \frac{m\pi y}{H}$

From or the possibility:

$$a_{nm} = \frac{4}{2H} \int_0^H \int_0^L f \sin \left(\frac{n\pi x}{L} \right) \sin \left(\frac{m\pi y}{H} \right) dx dy$$

As $t \rightarrow \infty \Rightarrow u \rightarrow 0$ (since all $\lambda > 0$)
 - the only equilibrium temperature with $u|_r = 0$.

(3)

$$\textcircled{c} \quad \frac{\partial u}{\partial x}(0, y, t) = \frac{\partial u}{\partial x}(L, y, t) = 0$$

$$u(x, 0, t) = 0, \quad u(x, H, t) = 0$$

$$u = \Phi(x, y) h(t)$$

$$\frac{d h}{d t} = -\lambda u h$$

$$\nabla^2 \Phi = -\lambda \Phi$$

Separ. of variables:

$$\Phi = f(x) g(y)$$

$$\Rightarrow \begin{cases} \frac{d^2 f}{dx^2} = -\mu f \\ f'(0) = f'(L) = 0 \end{cases}$$

$$\Rightarrow \mu = \left(\frac{n\pi}{L}\right)^2$$

Eigenfunktion:

$$f = \cos \frac{n\pi x}{L}, \quad n = 0, 1, \dots$$

$$\begin{cases} \frac{d^2 g}{dy^2} = -(\lambda - \mu) g \\ g(0) = g(H) = 0 \end{cases} \Rightarrow \lambda - \mu = \left(\frac{m\pi}{H}\right)^2, \quad m = 1, 2, \dots$$

$$\Rightarrow \lambda_{nm} = \left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2$$

Eigen function:

$$\phi = \cos \frac{n\pi x}{L} \sin \frac{m\pi y}{H}, \quad n = 0, 1, \dots$$

$$m = 1, 2, \dots$$

Superposition:

$$u = \sum_{n=0}^{\infty} \sum_{m=1}^{\infty} a_{nm} \cos \left(\frac{n\pi x}{L} \right) \sin \left(\frac{m\pi y}{H} \right) e^{-\lambda_{nm} k t}$$

$$\text{I.C.: } f = \sum_{n=0}^{\infty} \sum_{m=1}^{\infty} a_{nm} \cos \left(\frac{n\pi x}{L} \right) \sin \frac{m\pi y}{H}$$

By orthogonality:

$$a_{nm} = \frac{\int_0^H \int_0^L f \cos \frac{n\pi x}{L} \sin \frac{m\pi y}{H} dx dy}{\int_0^H \int_0^L \cos^2 \left(\frac{n\pi x}{L} \right) \sin^2 \left(\frac{m\pi y}{H} \right) dx dy}$$

here

$$\int_0^H \int_0^L \cos^2 \left(\frac{n\pi x}{L} \right) \sin^2 \left(\frac{m\pi y}{H} \right) dx dy = \begin{cases} \frac{1}{2} LH, & n=0 \\ \frac{1}{4} LH, & n \geq 1 \end{cases}$$

As $t \rightarrow \infty \Rightarrow u \rightarrow 0$ since zero BC for u

at $y=0$ and $y=H$ implies that $u=0$ for $t > 0$.

7.3.2d

5

$$u_t = \kappa (u_{xx} + u_{yy} + u_{zz})$$

$$\begin{aligned} 0 < x < L \\ 0 < y < H \\ 0 < z < W \end{aligned}$$

$$\begin{aligned} \frac{\partial u}{\partial x}(0, y, z, t) &= \frac{\partial u}{\partial x}(L, y, z, t) \\ &= \frac{\partial u}{\partial y}(x, 0, z, t) = \frac{\partial u}{\partial y}(x, H, z, t) = \\ &= \frac{\partial u}{\partial z}(x, y, 0, t) = \frac{\partial u}{\partial z}(x, y, W, t) = 0 \end{aligned}$$

$$u(x, y, z, 0) = \alpha(x, y, z)$$

$$u(x, y, z, t) = \phi(x, y, z) h(t)$$

$$\Rightarrow \frac{dh}{dt} = -\lambda \kappa h$$

$$\nabla^2 \phi = -\lambda \phi$$

Extra separation of variables:

$$\phi(x, y) = f(x)g(y)Q(z)$$

$$\begin{aligned} \nabla^2 \phi + \lambda \phi &= f_{xx}gQ + fg_{yy}Q + fgQ_{zz} + \lambda fgQ \\ &= 0 \end{aligned}$$

$$\Rightarrow \underbrace{\frac{f_{xx}}{f}}_{\mu} + \underbrace{\frac{g_{yy}}{g}}_{-\nu} + \underbrace{\frac{Q_{zz}}{Q} + \lambda}_{\mu + \nu} = 0$$

$$\Rightarrow \begin{cases} \frac{d^2 f}{dx^2} = -\mu f \\ f(0) = f(L) = 0 \end{cases}$$

$$\Rightarrow \mu = \left(\frac{n\pi}{L}\right)^2$$

$$f = \cos\left(\frac{n\pi x}{L}\right), \quad n=0, 1, \dots$$

$$\begin{cases} \frac{d^2 g}{dy^2} = -\nu g \\ g'(0) = g'(H) = 0 \end{cases} \Rightarrow$$

$$\nu = \left(\frac{m\pi}{H}\right)^2$$

$$g = \cos\left(\frac{m\pi y}{H}\right), \quad m=0, 1, \dots$$

$$\begin{cases} \frac{d^2 Q}{dz^2} = -(\lambda - \mu - \nu)Q \\ Q'(0) = Q'(W) = 0 \end{cases}$$

$$\Rightarrow \lambda - \mu - \nu = \left(\frac{l\pi}{W}\right)^2$$

$$Q = \cos\left(\frac{l\pi z}{W}\right), \quad l=0, 1, 2, \dots$$

$$\Rightarrow \lambda_{nml} = \left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2 + \left(\frac{l\pi}{W}\right)^2$$

$$\Phi = \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{H}\right) \cos\left(\frac{l\pi z}{W}\right)$$

$$n, m, l = 0, 1, 2, \dots$$

Superposition:

$$u = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} \sum_{l=0}^{\infty} a_{nml} \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{L}\right) \cos\left(\frac{l\pi z}{W}\right) e^{-\lambda_{nml} t}$$

IC: $f = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} \sum_{l=0}^{\infty} a_{nml} \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{L}\right) \cos\left(\frac{l\pi z}{W}\right)$

⇒ by orthogonality:

$$a_{nml} = \frac{\int_0^W \int_0^H \int_0^L \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{L}\right) \cos\left(\frac{l\pi z}{W}\right) f(x,y,z) dx dy dz}{\int_0^W \int_0^H \int_0^L \cos^2\left(\frac{n\pi x}{L}\right) \cos^2\left(\frac{m\pi y}{L}\right) \cos^2\left(\frac{l\pi z}{W}\right) dx dy dz}$$

As $t \rightarrow \infty$ all terms will decay but

$$A_{000} = \frac{1}{LHW} \int_0^L \int_0^H \int_0^W f(x,y,z) dx dy dz$$

$$\Rightarrow u(x,y,z,t) \Big|_{t \rightarrow \infty} = A_{000}$$

7.3.4. b

$$u_{tt} = c^2(u_{xx} + u_{yy})$$

$$0 < x < L, \quad 0 < y < H$$

IC: $u(x, y, 0) = 0 \quad \frac{\partial u}{\partial t}(x, y, 0) = \alpha(x, y)$

BC: $\frac{\partial u}{\partial x}(0, y, t) = \frac{\partial u}{\partial x}(L, y, t) = 0$

$\frac{\partial u}{\partial y}(x, 0, t) = \frac{\partial u}{\partial y}(x, H, t) = 0$

$u = \Phi(x, y, t) h(t)$

$\frac{d^2 h}{dt^2} = -\lambda c^2 h \Rightarrow h = c_1 \cos(\sqrt{\lambda} t) + c_2 \sin(\sqrt{\lambda} t)$
 $h(0) = 0 \Rightarrow c_1 = 0$ if $\lambda > 0$
 and $h = c_3 t$ if $\lambda = 0$

Separation of variables for Φ :

$\Phi = f(x) g(y)$

$\Rightarrow \begin{cases} \frac{d^2 f}{dx^2} = -\mu f \\ f'(0) = f'(L) = 0 \end{cases}$

$\Rightarrow \mu = \left(\frac{n\pi}{L}\right)^2$
 Eigenfunction

$f = \cos \frac{n\pi x}{L}, \quad n = 0, 1, \dots$

$$\begin{cases} \frac{d^2 g}{dy^2} = -(\lambda - \mu)g \\ g'(0) = g'(H) = 0 \end{cases}$$

$$\Rightarrow \lambda - \mu = \left(\frac{m\pi}{H}\right)^2$$

Eigen function

$$g = \cos\left(\frac{m\pi y}{H}\right), \quad m=0, 1, \dots$$

$$\Rightarrow \lambda_{nm} = \sqrt{\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2}$$

Superposition:

$$u = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} A_{nm} \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{H}\right) h_{nm}(t)$$

Bound of $h(t)=0$:

$$h_{nm}(t) = \begin{cases} t & n=m=0 \\ \sin(\omega_{nm} t) & \text{otherwise} \end{cases}$$

$$\text{here } \omega_{nm} = c \sqrt{\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2}$$

Sturm I.C: $\frac{\partial u}{\partial t} \Big|_{t=0} = \alpha(x, y) = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} A_{nm} \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{H}\right) \times h'_{nm}(0)$

$$\Rightarrow A_{nm} h'_{nm}(0) = \frac{\int_0^H \int_0^L f \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi y}{H}\right) dx dy}{\int_0^H \int_0^L \cos^2\left(\frac{n\pi x}{L}\right) \cos^2\left(\frac{m\pi y}{H}\right) dx dy}$$

Note that $h'_{nm}(0) = 1$ for $n=m=0$

but otherwise $h'_{nm}(0) = \omega_{nm}$

7.3.7c

$$u_{xx} + u_{yy} + u_{zz} = 0$$

$$0 < x < L, 0 < y < W, 0 < z < H$$

$$\frac{\partial u}{\partial x}(0, y, z) = 0$$

$$\frac{\partial u}{\partial x}(L, y, z) = \alpha(x, y, z)$$

$$\frac{\partial u}{\partial y}(x, 0, z) = 0$$

$$\frac{\partial u}{\partial y}(x, W, z) = 0$$

$$\frac{\partial u}{\partial z}(x, y, 0) = \frac{\partial u}{\partial z}(x, y, H) = 0$$

$$u = h(x) \Phi(y, z) \Rightarrow$$

(10)
(11)

$$\frac{d^2 h}{dx^2} = \beta h \quad \text{and}$$

$$\nabla^2 \Phi + \beta \Phi = 0$$

$$\Rightarrow h = \cosh(\sqrt{\beta} x)$$

$$\text{since } h'(0) = 0$$

$$\frac{\partial \Phi}{\partial y}(0, z) = \frac{\partial \Phi}{\partial y}(w, z) = 0$$

$$\frac{\partial \Phi}{\partial z}(y, 0) = \frac{\partial \Phi}{\partial z}(y, H) = 0$$

$$\text{sol. is: } \phi = f(y) g(z)$$

From (7.3.7. (b)) :

$$\begin{cases} \frac{d^2 f}{dy^2} = -\mu f \\ f'(0) = f'(w) = 0 \end{cases}$$

$$\mu = \left(\frac{n\pi}{w} \right)^2$$

Eigenfunction:

$$f = \cos \frac{n\pi y}{w}, \quad n=0, 1, \dots$$

$$\begin{cases} \frac{d^2 g}{dz^2} = -(\lambda - \mu) g \\ g'(0) = g'(H) = 0 \end{cases}$$

$$\Rightarrow \lambda - \mu = \left(\frac{m\pi}{H} \right)^2$$

Eigenfunction

$$g = \cos \left(\frac{m\pi z}{H} \right), \quad m=0, 1, \dots$$

$$\beta_{nm} = \left(\frac{n\pi}{w} \right)^2 + \left(\frac{m\pi}{H} \right)^2$$

Superposition:

(12)

$$u(x, y, z) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} A_{nm} \cos\left(\frac{n\pi y}{W}\right) \cos\left(\frac{m\pi z}{H}\right) \cosh(\beta_{nm} x)$$

Non homogeneous boundary condition:

$$\frac{\partial u}{\partial x} \Big|_{x=L} = \alpha(y, z) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} A_{nm} \beta_{nm} \cancel{\cos\left(\frac{n\pi y}{W}\right)} \cancel{\cos\left(\frac{m\pi z}{H}\right)} \\ \times \sinh(\beta_{nm} L) \cos\left(\frac{n\pi y}{W}\right) \cos\left(\frac{m\pi z}{H}\right)$$

At

\Rightarrow orthogonality:

$$A_{nm} \sinh(\beta_{nm} L) = \frac{\int_0^H \int_0^W \alpha(y, z) \cos\left(\frac{n\pi y}{W}\right) \cos\left(\frac{m\pi z}{H}\right) dy dz}{\int_0^H \int_0^W \cos^2 \frac{n\pi y}{W} \cos^2 \frac{m\pi z}{H} dy dz}$$

At $n=m=0$: $\sinh(0) = 0$

$$\Rightarrow \text{require } \int_0^H \int_0^W \alpha(y, z) dy dz = 0$$

$\Rightarrow A_{00}$ - arbitrary.

7.5.6

$$I = \iint_R [u L(v) - v L(u)] dx dy$$

$$L = \nabla^2 + q(x, y)$$

$$L(u) = \nabla^2 u + q u$$

$$L(v) = \nabla^2 v + q v$$

$$\Rightarrow u L(v) - v L(u) = u \nabla^2 v + \cancel{q u v} - v \nabla^2 u - \cancel{q v u}$$

$$= u \nabla^2 v - v \nabla^2 u$$

$$I = \iint_R [u L(v) - v L(u)] dx dy = \iint_R (u \nabla^2 v - v \nabla^2 u) dx dy$$

$$= \oint_{\partial R} (u \vec{\nabla} v - v \vec{\nabla} u) \cdot \vec{n} ds$$

①

boundary of R