

# Two weight estimates, Clark measures and rank one perturbations of unitary operators

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- WLOG:  $b$  is cyclic, so  $U = M_\xi$  in  $L^2(\mu)$ ,  $\mu(\mathbb{T}) = 1$ ;  $b \equiv \mathbf{1}$ , therefore  $b_1(\xi) = \bar{\xi}$ .

# Spectral theorem

- If  $U$  is cyclic, i.e. if for some  $b \in \mathcal{H}$  we have  $\text{span}\{U^n b : n \in \mathbb{Z}\} = \mathcal{H}$ , then  $U = M_\xi$  in  $L^2(\mu)$ . Vector  $b$  is called *cyclic* vector for  $U$ .
- Measure  $\mu$  is not unique, but we can pick a measure  $\mu = \mu_b$  associated to  $b$ ,

$$(U^n b, b) = \int_{\mathbb{T}} \xi^n d\mu_b(\xi) \quad \forall n \in \mathbb{Z},$$

the measure  $\mu_b$  is uniquely defined.

- If  $\Phi : \mathcal{H} \rightarrow L^2(\mu)$ ,  $\mu = \mu_b$  is the unitary operator such that  $U = \Phi^{-1} M_\xi \Phi$ , then  $\Phi b = \mathbf{1}$ .

# Back to rank 1 perturbations

- $U_\gamma = U + (\gamma - 1)bb_1^*$ ,  $\gamma \in \mathbb{T}$ ,  $b$  cyclic,  $b_1 = U^*b$ .
- Let  $b$  is a cyclic vector for  $U$ . It is not hard to show that then  $b$  is cyclic for all  $U_\gamma$ .
- We consider spectral measures  $\mu_\gamma = \mu_{\gamma,b}$  for operators  $U_\gamma$ ,  
In this case  $\Phi_\gamma b \equiv \mathbf{1}$  for all  $\gamma \in \mathbb{T}$ .  
Since  $\|b\| = 1$  all measures  $\mu_\gamma$  are probability measures.
- The measures  $\mu_\gamma$  are called Clark measures.
- So let  $U = U_1$  be  $M_\xi$  in  $L^2(\mu)$ ,  $\mu = \mu_1$ .

Then

$$b \equiv 1, \quad b_1 = U^*b \equiv \bar{\xi}.$$

**Goal:** Want to describe unitary operators intertwining  $U_\gamma$  and its model:  
unitary  $\Phi_\gamma : L^2(\mu) \rightarrow L^2(\mu_\gamma)$ ,

$$U_\gamma \Phi_\gamma = M_z \Phi, \quad U_\gamma b \equiv \mathbf{1}.$$

# Relations for Clark measures

- The measure  $\mu_\gamma$  is defined

$$((U_\gamma - zI)^{-1}b, b) = \int_{\mathbb{T}} \frac{d\mu_\gamma(\xi)}{z - \xi} \quad \forall z \notin \mathbb{T}.$$

- $U_\gamma$  is a rank 1 perturbation of  $U_1 = M_\xi$  in  $L^2(\mu)$ . Using the formula

$$(I - ac^*)^{-1} = I + \frac{1}{d}ac^*, \quad d = 1 - (c, a),$$

one can compute the resolvent and find relations between  $\mu_\gamma$ .



# Relations for Clark measures

- Namely, let

$$T\mu(z) = \int_{\mathbb{T}} \frac{d\mu(\xi)}{1 - \bar{\xi}z}.$$

- Define  $\theta(z)$  by

$$1 - \theta(z) = \frac{1}{T\mu(z)}$$

- Then

$$1 - \bar{\gamma}\theta(z) = \frac{1}{T\mu_{\gamma}(z)}$$

- Can rewrite

$$T\mu_{\gamma}(z) = \frac{T\mu(z)}{1 - (1 - \gamma)T\mu(z)}.$$

# Relations for Clark measures

Function  $\theta$  was introduced for a reason.

- One can show that  $\theta \in H^\infty$ ,  $\|\theta\|_\infty \leq 1$ .
- Therefore the function  $F_\gamma = \frac{1 + \bar{\gamma}\theta}{1 - \bar{\gamma}\theta}$  has positive real part.
- The measures  $\mu_\gamma$  are the measures whose Poisson extension give  $\operatorname{Re} F_\gamma$ .

# Pretend to be physicists

Let  $\Phi$  be an integral operator with kernel  $K(z, \xi)$

$$\Phi f(z) = \int_{\mathbb{T}} K(z, \xi) f(\xi) d\mu(\xi).$$

Want:

$$\Phi_{\gamma}(M_{\xi} + (\gamma - 1)bb_1^*) = M_z\Phi_{\gamma}.$$

Recall that  $b \equiv \mathbf{1}$ ,  $\Phi_{\gamma}b \equiv \mathbf{1}$  (in  $L^2(\mu_{\gamma})$ ),  $b_1 = U^*b \equiv \bar{\xi}$  (in  $L^2(\mu)$ ), so  $\Phi_{\gamma}bb_1^*$  is an integral operator with kernel  $1 \cdot \xi$ .

$$K(z, \xi)\xi + (\gamma - 1)\xi = zK(z, \xi).$$

Solving for  $K$  we get

$$K(z, \xi) = (1 - \gamma) \frac{\xi}{\xi - z} = (1 - \gamma) \frac{1}{1 - \bar{\xi}z}$$

# Acting as boring mathematicians

Theorem (C. Liaw, S. Treil)

*Under the above assumptions*

$$\Phi_\gamma f(z) = f(z) + (1 - \gamma) \int_{\mathbb{T}} \frac{f(\xi) - f(z)}{1 - \bar{\xi}z} d\mu(\xi).$$

*for all  $f \in C^1(\mathbb{T})$ .*

Is it a SIO with kernel  $1/(1 - \bar{\xi}z)$ ?

- Will show next lecture that the operators  $T_r$  with kernels  $1/(1 - r\bar{\xi}z)$ ,  $\xi, z \in \mathbb{T}$  are uniformly (in  $r \neq 1$ ) bounded (it also follows from results of V. Kapustin).
- Follows from results of A. Poltoratskii that the boundary values  $T_\pm f$  (from inside and outside) of the Cauchy integral

$$Tf(z) = \int_{\mathbb{T}} \frac{f(\xi)}{1 - \bar{\xi}z} d\mu(\xi)$$

exist  $\mu_\gamma$ -a.e.

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- That implies the operators  $T_+ : L^2(\mu) \rightarrow L^2(\mu_\gamma)$  are bounded and

$$\Phi_\gamma f = (\mathbf{1} - (1 - \gamma)T_\pm \mathbf{1})f + (1 - \gamma)T_\pm f$$

From the previous slide:

$$\Phi_\gamma f = (\mathbf{1} - (1 - \gamma)T_\pm \mathbf{1})f + (1 - \gamma)T_\pm f$$

It is known (Aronszajn–Donoghue) that  $\mu_s \perp (\mu_\gamma)_s$ , so

$$T_\pm \mathbf{1} = \frac{1}{1 - \gamma} \quad (\mu_\gamma)_s\text{-a.e.},$$

which agrees with known results.

# Proof of the representation formula

$$\Phi_\gamma f(z) = f(z) + (1 - \gamma) \int_{\mathbb{T}} \frac{f(\xi) - f(z)}{1 - \bar{\xi}z} d\mu(\xi).$$

Idea of the proof:

$$U_\gamma = M_\xi + (\gamma - 1)bb_1^*, \quad b = \mathbf{1}, \quad b_1 \equiv \xi.$$

and

$$\Phi_\gamma(M_\xi + (\gamma - 1)bb_1^*) = M_z\Phi_\gamma.$$

can be rewritten as

$$\Phi_\gamma M_\xi = M_z\Phi_\gamma + (1 - \gamma)(\Phi_\gamma b)b_1^*.$$

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Right multiplying (\*) by  $M_\xi$  and using (\*) we get

$$\begin{aligned} \Phi_\gamma M_\xi^2 &= M_z \Phi_\gamma M_\xi + (1 - \gamma)(\Phi_\gamma b)b_1^* M_\xi \\ &= M_z^2 \Phi_\gamma + (1 - \gamma)(M_z(\Phi_\gamma b)b_1^* + (\Phi_\gamma b)b_1^* M_\xi). \end{aligned}$$

Iterating and using  $\Phi_\gamma b = \mathbf{1}$  we get

$$\Phi_\gamma M_\xi^n = M_z^n \Phi_\gamma + (1 - \gamma) \sum_{k=0}^{n-1} (M_z^k \mathbf{1}) b_1^* M_\xi^{n-k-1}.$$

Applying this identity to  $b = \mathbf{1}$  we get for  $f(\xi) \equiv \xi^n$ ,  $n \geq 0$

$$\begin{aligned} \Phi_\gamma f(z) &= z^n + (1 - \gamma) \sum_{k=0}^{n-1} z^k \int_{\mathbb{T}} \xi^{n-k} d\mu(\xi) \\ &= z^n + (1 - \gamma) \int_{\mathbb{T}} \frac{\xi^n - z^n}{1 - \bar{\xi}z} d\mu(\xi). \end{aligned}$$



Applying the identity

$$\Phi_\gamma M_\xi^n = M_z^n \Phi_\gamma + (1 - \gamma) \sum_{k=0}^{n-1} (M_z^k \mathbf{1}) b_1^* M_\xi^{n-k-1}.$$

to  $f(\xi) \equiv \bar{\xi}^n$  and multiplying the result by  $\bar{z}^n$  we get

$$\begin{aligned} \bar{z}^n &= (\Phi_\gamma f)(z) + (1 - \gamma) \bar{z}^n \sum_{k=0}^{n-1} z^k \int_{\mathbb{T}} \xi^{n-k} \bar{\xi}^n d\mu(\xi) \\ &= (\Phi_\gamma f)(z) + \int_{\mathbb{T}} \bar{z}^n \bar{\xi}^n \frac{\xi^n - z^n}{1 - \bar{\xi}z} d\mu(\xi) \end{aligned}$$

# Rigidity theorem

## Theorem (Rigidity Theorem)

*Let a probability measure  $\mu$  on  $\mathbb{T}$  be supported on at least two distinct points. Let  $\gamma \in \mathbb{T} \setminus \{1\}$ , and let  $\mathcal{V}f$  be defined for  $C^1$  functions  $f$*

$$\mathcal{V}f(z) = f(z) + (1 - \gamma) \int_{\mathbb{T}} \frac{f(\xi) - f(z)}{1 - \bar{\xi}z} d\mu(\xi)$$

*Assume  $\mathcal{V}$  extends to a bounded operator from  $L^2(\mu)$  to  $L^2(\nu)$  and assume  $\ker \mathcal{V} = \{0\}$ .*

*Then  $\exists h > 0$  such that  $1/h \in L^\infty(\nu)$ , and  $M_h \mathcal{V}$  is a unitary operator from  $L^2(\mu) \rightarrow L^2(\nu)$  (equivalently, that  $\mathcal{V} : L^2(d\mu) \rightarrow L^2(|h|^2 d\nu)$  is unitary).*

*Moreover, the measure  $|h|^2 \nu$  is exactly the Clark measure  $\mu_\gamma$ , and  $\mathcal{V}$  treated as the operator  $L^2(\mu) \rightarrow L^2(\mu_\gamma)$  is exactly the operator  $\Phi_\gamma$ .*

The main idea of the proof: similar unitary operators are unitarily equivalent