

GAUGE FIELDS, MATTER FIELDS, GRAV. FIELDS

Alg. situation

$k = \bar{k}$ alg of field

M non sing var / k , G alg gr / k , $G \subset GL_n$, $\mathfrak{g} = \text{Lie}(G)$ closed

V fin dim k -lin sp., $\rho: G \rightarrow GL(V)$ mor, $\rho^! = \text{ad} \rho: \mathfrak{g} \rightarrow \mathfrak{gl}(V)$
 $\Omega^i(M) = i\text{-forms on } M$, $G(M) = \text{Hom}_k(M, G)$.

C^∞ situation

(similar) will wk in alg. sit for now but also in C^∞ sit when mks surx.

Note ~~in alg~~

$G \times V \rightarrow V$, $G \times \mathfrak{g} \rightarrow \mathfrak{g}$ & map $\mathfrak{g} \times V \xrightarrow{\rho^! = \text{ad} \rho} V$ is G -equiv

Let $U(\mathfrak{g}) = \hat{U}$ be univ. envelop. algebra. Define

- (D) $[,] = \wedge_{\mathfrak{g}}: (\Omega^i(M) \otimes \mathfrak{g}) \times (\Omega^j(M) \otimes \mathfrak{g}) \rightarrow \Omega^{i+j}(M) \otimes \mathfrak{g}$
 $\wedge_{\hat{U}}: (\Omega^i(M) \otimes U(\mathfrak{g})) \times (\Omega^j(M) \otimes U(\mathfrak{g})) \rightarrow \Omega^{i+j}(M) \otimes U(\mathfrak{g})$
 $\wedge_{\rho}: (\Omega^i(M) \otimes \mathfrak{g}) \times (\Omega^j(M) \otimes V) \rightarrow \Omega^{i+j}(M) \otimes V$
 (R) If $A = \sum \alpha_i \otimes \xi_i$, $B = \sum \beta_j \otimes \zeta_j \Rightarrow A \wedge_{\rho} B = \sum (\alpha_i \wedge \beta_j) \otimes [\xi_i, \zeta_j]$, $[A, B] = \sum (\alpha_i \wedge \beta_j) \otimes [\xi_i, \zeta_j]$
 (D) $d: \Omega^i(M) \otimes \mathfrak{g} \rightarrow \Omega^{i+1}(M) \otimes \mathfrak{g}$, $d(A \wedge_{\rho} B) = dA \wedge_{\rho} B + (-1)^i A \wedge_{\rho} dB$
 $\& \text{ sim for } U(\mathfrak{g}) \& V.$

(D) Elts of $\Omega^1(M) \otimes \mathfrak{g}$ called gauge fields

Group $G(M)$ called gauge group
 Elts of $\Omega^1(M) \otimes V$ called matter fields.

$\Omega^1(M) \otimes \mathfrak{gl}_n(k)$

$GL_n(\mathcal{O}(M)) \quad \mathfrak{gl}_n(\Omega^1(M))$

(P) (Log. derivative) The map $G(M) \rightarrow \Omega^1(M) \otimes \mathfrak{gl}_n(k)$, $\gamma \mapsto \gamma^{-1} d\gamma$ takes values in $\Omega^1(M) \otimes \mathfrak{g}$.

Pf Show $\forall \partial \in \text{Der}_k(\mathcal{O}(M), \mathcal{O}(M))$, $\gamma^{-1} \partial \gamma \in \mathfrak{gl}_n(\mathcal{O}(M))$. Say $G = Z(F_\alpha)$.
 So $F_\alpha(\gamma) = 0$ so $(1 + \epsilon \partial)(F_\alpha(\gamma)) = 0$ in $k[\epsilon]$.

" $\& 1 + \epsilon \partial$ ring homo

$$F_\alpha((1 + \epsilon \partial)(\gamma))$$

"

$$F_\alpha(\gamma + \epsilon \partial \gamma)$$

$$\Rightarrow \gamma + \epsilon \partial \gamma \in G(k[\epsilon]) \Rightarrow \gamma^{-1} \in G(k[\epsilon]) \Rightarrow \gamma^{-1} \partial \gamma \in \mathfrak{gl}_n(k[\epsilon])$$

" $\& \gamma^{-1} \partial \gamma \in \mathfrak{g}$
 $\& \gamma^{-1} \partial \gamma \in \mathfrak{gl}_n(k)$

(Constr) Let $G(M) \times (\Omega^1(M) \otimes \mathfrak{g}) \rightarrow \Omega^1(M) \otimes \mathfrak{g}$

be action $(\gamma^{-1}, A) \mapsto A^{\gamma^{-1}} := \gamma^{-1} A \gamma + \gamma^{-1} d\gamma$ called gauge action

So gauge gr acts on gauge fields. $(\text{ad } \gamma^{-1})(A) = \rho^!(\gamma^{-1}) A$

$\gamma^{-1} \partial \gamma \in \mathfrak{g}(k)$
 for each pt in M .

(Rem) Also have action $G(M) \times (\Omega^i(M) \otimes \mathfrak{g}) \rightarrow \Omega^i(M) \otimes \mathfrak{g}$ by ad , $(\gamma, B) \mapsto \gamma B \gamma^{-1} = \rho^!(\gamma) B$.

(Def) Curvature of $A \in \Omega^1(M) \otimes \mathfrak{g}$ is $F = F_A := dA + A \wedge A$

$$(A = \sum \alpha_i \otimes \xi_i \Rightarrow [A, A] = \sum (\alpha_i \wedge \alpha_j) (\xi_i \otimes \xi_j - \xi_j \otimes \xi_i) = 2A \wedge A) = dA + \frac{1}{2} [A, A] \in \Omega^2(M) \otimes \mathfrak{g}$$

(Def) For $A \in \Omega^1(M) \otimes \mathfrak{g}$ def covariant exterior differentiation $D = D_A$

$$D\Phi = d\Phi + A \wedge \Phi, \quad D: \Omega^i(M) \otimes V \rightarrow \Omega^{i+1}(M) \otimes V$$

* acting trivially on $\Omega^i(M)$ & simil have action $G(M) \times (\Omega^i(M) \otimes V) \rightarrow \Omega^i(M) \otimes V$ ind by ρ .

For $\Phi \in \Omega^i(M) \otimes V$ -2- $\Omega^i(M) \otimes V \rightarrow \Omega^{i+2}(M) \otimes V$

① $D_A^2 \Phi = F_A \wedge \Phi$ (in partic D_A^2 is $\mathcal{O}(M)$ -linear)

Pf. $D^2 \Phi = d(d\Phi + A \wedge \Phi) + A \wedge (d\Phi + A \wedge \Phi) \quad (A \wedge A) \wedge \Phi$
 $= 0 + dA \wedge \Phi - \cancel{A \wedge d\Phi} + \cancel{A \wedge d\Phi} + A \wedge (A \wedge \Phi) = F_A \wedge \Phi.$

② For $\Phi \in \Omega^i(M) \otimes V$ have $D_A(\gamma \Phi) = \gamma D_{A^{\gamma^{-1}}} \Phi$ (covariance of D)

Pf. $\gamma^{-1} D_A \gamma \Phi = \gamma^{-1} [d(\gamma \Phi) + A \wedge \gamma \Phi]$
 $= \gamma^{-1} [d\gamma \wedge \Phi + \gamma d\Phi] + \gamma^{-1} [A \wedge \gamma \Phi]$
 $= \gamma^{-1} d\gamma \wedge \Phi + d\Phi + \gamma^{-1} A \gamma \wedge \gamma^{-1} \gamma \Phi$
 $= d\Phi + (\underbrace{\gamma^{-1} d\gamma + \gamma^{-1} A \gamma}_{A^{\gamma^{-1}}}) \wedge \Phi = D_{A^{\gamma^{-1}}} \Phi.$

(of $\gamma: V \rightarrow V$ is G -equiv)

③ $F_{A^{\gamma^{-1}}} = \gamma^{-1} F_A \gamma$ (covariance of F)

Pf. Enough (check!) $F_{A^{\gamma^{-1}}} \Phi = \gamma^{-1} F_A \gamma \wedge \Phi$ for all $\gamma \in \Phi$.

Now $F_{A^{\gamma^{-1}}} \wedge \Phi = D_{A^{\gamma^{-1}}}^2 \Phi = D_{A^{\gamma^{-1}}} D_{A^{\gamma^{-1}}} \Phi = \gamma^{-1} D_A (\gamma \gamma^{-1} D_A (\gamma \Phi))$
 $= \gamma^{-1} D_A^2 (\gamma \Phi) = \gamma^{-1} (F_A \wedge \gamma \Phi) = (\gamma^{-1} F_A \gamma) \wedge \Phi \quad \square.$

④ (Bianchi id) Look @ case $V = \mathfrak{g}$, $\rho = \text{ad}$. So $\begin{cases} A \wedge B = [A, B] \\ = A \wedge B - B \wedge A \end{cases}$
 for $A \in \Omega^1(M) \otimes \mathfrak{g}$, $B \in \Omega^2(M) \otimes \mathfrak{g}$. Then $\begin{cases} dF_A = [F_A, A] \\ D_A F_A = 0 \end{cases}$

Pf. $dF_A = d(dA + A \wedge A) = \underbrace{d^2 A}_0 + \underbrace{dA \wedge A}_0 - \underbrace{A \wedge dA}_0$
 $= (F - A \wedge A) \wedge A - A \wedge (F - A \wedge A)$
 $= F \wedge A - A \wedge F = [F, A] = -[A, F] = -A \wedge F. \quad \text{So } D_A F_A = 0$
 $(\gamma, (\Phi, A)) \mapsto (\gamma \Phi, A^\gamma)$

⑤ $G(M)$ acts on $[\mathcal{O}(M) \otimes V] \times [\Omega^1(M) \otimes \mathfrak{g}]$ (ρ on 1st factor, gauge action on 2nd)

matter fields $\mathcal{F}_{\text{mat}}$ gauge fields $\mathcal{F}_{\text{gauge}}$

"diff-als"
 A fcn $L: \mathcal{F}_{\text{mat}} \times \mathcal{F}_{\text{gauge}} \rightarrow \Omega^d(M)$ is called a Lagrangian. L called gauge inv if L const on $G(M)$ -orbits. ($d = \dim M$)

Σ is induced by sheaf isom $\Omega_M^{d-1} \cong \Omega_M^{d-1}$

Def A k -lin op $*$: $\bigoplus \Omega^i(M) \rightarrow \bigoplus \Omega^i(M)$ called a Hodge op if bij $\Sigma * \Omega^i(M) = \Omega^{d-i}(M)$.
It induces $G(M)$ -equiv $\omega_{ij}^{d-i=0} * : \Omega^i(M) \otimes_{\mathcal{O}_M} \mathcal{O}_M \rightarrow \Omega^{d-i}(M) \otimes_{\mathcal{O}_M} \mathcal{O}_M$ where $G(M)$ acts on $\Omega^i(M) \otimes_{\mathcal{O}_M} \mathcal{O}_M$ via ad.

EG (examples of gauge inv Lagr's) $\dim M = d = 4$

1) ~~Let~~ $L_1(\Phi, A) := \text{tr} (F_A \wedge * F_A)$ (where $\text{tr} : \Omega^4(M) \otimes_{\mathcal{O}_M} \mathcal{O}_M \rightarrow \Omega^4(M)$, $\text{tr} : \mathcal{O}_M \subset \text{gl}_n(k) \rightarrow k$)

Indeed $L_1(\gamma^{-1}\Phi, A^{\gamma^{-1}}) = \text{tr} (F_{A^{\gamma^{-1}}} \wedge * F_{A^{\gamma^{-1}}}) = \text{tr} (\gamma^{-1} F_A \gamma \wedge * \gamma^{-1} F_A \gamma) = \text{tr} (\gamma^{-1} F_A \gamma \wedge \gamma^{-1} * F_A \gamma) = \text{tr} (\gamma^{-1} F_A \wedge * F_A \gamma) = \text{tr} (F_A \wedge * F_A)$

2) $L_2(\Phi, A) = \Phi^t \wedge * \Phi$ (where $V = \text{Mat}_{n,1} \ni \Phi \in \text{Mat}_{n,1}$ is the transpos w/ $G \subset O(n)$ acting on $V \cong V^t$ naturally).

Indeed $(\gamma \Phi)^t \wedge * \gamma \Phi = \Phi^t \gamma^t \wedge * \gamma \Phi = \Phi^t \wedge * \Phi$

3) $L_3(\Phi, A) = (D_A \Phi)^t \wedge * D_A \Phi$ (again $G \subset O(n)$, $V = \text{Mat}_{n,1}$ etc as in 2))

Indeed $(D_{A^{\gamma^{-1}}}(\gamma \Phi))^t \wedge * D_{A^{\gamma^{-1}}}(\gamma \Phi) = (D_A \Phi)^t \gamma^t \wedge * \gamma D_A \Phi = (D_A \Phi)^t \wedge * D_A \Phi$

4) $L_{YM} = L_1 + L_2 + L_3$ (Yang Mills Lagrangian, in the $U(n)$ version)

D Let E be a vect bdl on M . A connection on E is a k -lin map $\nabla : E \rightarrow E \otimes \Omega_M^1$ s.t. $\nabla(fe) = df \otimes e + f \nabla e$. $\exists ! \nabla : E \otimes \Omega_M^0 \rightarrow E \otimes \Omega_M^1$ s.t. ∇ is the dd ∇ on E & $\nabla(e \otimes w) = \nabla e \wedge w + e \otimes dw$. This has also prop $\nabla(\eta \wedge w) = \nabla \eta \wedge w + \eta \wedge dw$.

It foll that $\nabla^2(fe) = \nabla(e \otimes df + \nabla e \otimes f) = \nabla e \otimes df + \nabla^2 e \otimes f - \nabla e \otimes df$ so $\nabla^2 : E \rightarrow E \otimes \Omega^2$ is $\mathcal{O}$ -linear so it def an elt of $H^0((\text{End } E) \otimes \Omega^2) = \text{gl}_n(k) \otimes_k \Omega^2(M)$ called F_{∇} (curv of ∇)

P In not's above, assume E trivial. Then 1) $\mathcal{O}(M) \otimes V = \mathcal{O}(E) = \{\text{sect's of } E\}$ matter fields

2) $\frac{\Omega^1(M) \otimes \text{gl}_n(k)}{\text{GL}_n(M)} \cong \{ \text{connections on } E \}$ & curvatures correspond

Pf If ∇ conn $\Sigma(e_i)$ basis write $\nabla e_i = \Sigma e_j \otimes \omega_{ij}$ & attach (ω_{ij}) mod gauge eq.

$\exists$ obvious version with w/ $O(n)$ repl by $U(n)$ in which $()^t$ repl by $()^*$.

This suggests to see connections on (non-triv) vect bdl's (or bdl's w/ other gr's than GL_n) as (generalisations of) gauge fields. Sections in such bdl's are (generalized) matter fields

Set $T = T_M, \Omega = \Omega^1_M$

① A gravit field is an elt $g \in (\Omega \otimes \Omega)_0 \subset \Omega \otimes \Omega$ (also called metric) where "index 0" means "non-deg".

If ω^α basis of Ω , $g = g_{\alpha\beta} \omega^\alpha \otimes \omega^\beta$ (Einst not'n, $g_{\alpha\beta} = g_{\beta\alpha}$) non-deg means $\det g := \det g_{\alpha\beta} \neq 0$

② A connection on T is (cf gen case above) a k -lin map $\nabla: T \rightarrow T \otimes \Omega$ If ω^α as above & ∂_α basis of T dual to ω^α (ie $\partial_\alpha: \mathcal{O} \rightarrow \mathcal{O}$ k -der, $\{\partial_\alpha, \omega^\beta\} = \delta^\beta_\alpha$) then $\nabla \partial_\mu = \Gamma^\nu_{\mu\sigma} \partial_\nu \otimes \omega^\sigma$, $df = \partial_\mu f \cdot \omega^\mu$

"Christoffel symbols"

$$\begin{aligned} \text{so } \nabla(V^\mu \partial_\mu) &= \partial_\mu \otimes dV^\mu + \nabla \partial_\mu \cdot V^\mu = \partial_\mu \otimes \partial_\sigma V^\mu \cdot \omega^\sigma + \Gamma^\nu_{\mu\sigma} \partial_\nu \otimes \omega^\sigma \\ &= [\partial_\sigma V^\mu + \Gamma^\mu_{\sigma\mu} V^\mu] \partial_\nu \otimes \omega^\sigma \end{aligned}$$

def via $E \otimes \Omega \otimes F \cong E \otimes F \otimes \Omega$

$$\nabla(e \otimes f) = \nabla e \otimes f + e \otimes \nabla f$$

③ Have natural $\otimes$ & dual of vect bdl's w/ conn's s.t. In partic ∇ on T induces ∇ on Ω & on $T \otimes \dots \otimes T \otimes \Omega \otimes \dots \otimes \Omega = T^{\otimes p} \otimes \Omega^{\otimes q}$ For $\tau \in T^{\otimes p} \otimes \Omega^{\otimes q}$ write $\tau = \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \partial_\alpha \otimes \partial_\beta \otimes \dots \otimes \omega^{\alpha'} \otimes \omega^{\beta'} \otimes \dots$

Image of τ via $T^{\otimes p} \otimes \Omega^{\otimes q} \rightarrow T^{\otimes p-1} \otimes \Omega^{\otimes q-1}$ induced by $T \otimes \Omega \rightarrow \mathcal{O}$ (on 1st place) is $(\tau^{\alpha\beta\dots}_{\alpha'\beta'\dots}) \partial_\beta \otimes \partial_\gamma \otimes \dots \otimes \omega^{\beta'} \otimes \omega^{\gamma'} \otimes \dots$

Note $\nabla \omega^\nu = -\Gamma^\nu_{\mu\sigma} \omega^\mu \otimes \omega^\sigma$ (b/c $0 = d\langle \partial_\mu, \omega^\nu \rangle = \langle \nabla \partial_\mu, \omega^\nu \rangle + \langle \partial_\mu, \nabla \omega^\nu \rangle$)

More generally $\nabla(\tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \partial_\alpha \otimes \dots \otimes \omega^{\alpha'} \otimes \dots) = \leftarrow \partial_\alpha \otimes \dots \otimes \omega^{\alpha'} \otimes \dots \otimes d\tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} + \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \nabla \partial_\alpha \otimes \dots \otimes \omega^{\alpha'} \otimes \dots + \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \partial_\alpha \otimes \dots \otimes \nabla \omega^{\alpha'} \otimes \dots$

(*)

$$\begin{aligned} &= \partial_\sigma \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \cdot \partial_\alpha \otimes \dots \otimes \omega^{\alpha'} \otimes \dots \otimes \omega^\sigma + \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \Gamma^\nu_{\alpha\sigma} \partial_\nu \otimes \dots \otimes \omega^{\alpha'} \otimes \dots \otimes \omega^\sigma \\ &= (\partial_\sigma \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} + \Gamma^\alpha_{\nu\sigma} \tau^{\nu\beta\dots}_{\alpha'\beta'\dots} + \dots - \Gamma^{\mu'}_{\alpha'\sigma} \tau^{\alpha\beta\dots}_{\mu'\beta'\dots} - \dots) \partial_\alpha \otimes \dots \otimes \omega^{\alpha'} \otimes \dots \otimes \omega^\sigma \\ &\equiv: \nabla_\sigma \tau^{\alpha\beta\dots}_{\alpha'\beta'\dots} \end{aligned}$$

The "Riemann tensor"

(R) ∇_μ sat "Leibn rule" in the sense that

$$\nabla_\mu (V^{\alpha\beta\dots} W^{\gamma\delta\dots}) = (\nabla_\mu V^{\alpha\beta\dots}) W^{\gamma\delta\dots} + V^{\alpha\beta\dots} \nabla_\mu W^{\gamma\delta\dots}$$

& similarly if some indices contracted (e.g. $\alpha = \gamma$)

[Follows from fact that ∇ sat Leibniz wrt to $\otimes$ & $\langle \rangle$].

(R) Write $\nabla^2 \omega^\sigma_\mu = R^\alpha_{\mu\rho\sigma} \partial_\alpha \omega^\rho_\lambda \omega^\lambda$ & compute $R^\alpha_{\mu\rho\sigma}$ in terms of Γ

$$\begin{aligned} \nabla^2 \omega^\sigma_\mu &= \nabla(\Gamma^\nu_{\mu\sigma} \partial_\nu \omega^\sigma) = \partial_\nu \omega^\sigma \wedge d\Gamma^\nu_{\mu\sigma} + \Gamma^\nu_{\mu\sigma} \nabla \partial_\nu \omega^\sigma + \Gamma^\nu_{\mu\sigma} \partial_\nu \omega^\sigma \wedge \nabla \omega^\sigma \\ &= \partial_\nu \omega^\sigma \wedge \partial_\sigma \Gamma^\nu_{\mu\sigma} \cdot \omega^\sigma + \Gamma^\nu_{\mu\sigma} \Gamma^k_{kj} \partial_\nu \omega^\sigma \wedge \omega^j - \Gamma^\nu_{\mu k} \Gamma^k_{\sigma j} \partial_\nu \omega^\sigma \wedge \omega^j \\ &= (\partial_j \Gamma^\nu_{\mu\sigma} + \Gamma^\nu_{\mu\sigma} \Gamma^k_{kj} - \Gamma^\nu_{\mu k} \Gamma^k_{\sigma j}) \partial_\nu \omega^\sigma \wedge \omega^j \end{aligned}$$

So $F^\nu_\mu = \partial_j \Gamma^\nu_{\mu\sigma} + \Gamma^\nu_{\mu\sigma} \Gamma^k_{kj} - \Gamma^\nu_{\mu k} \Gamma^k_{\sigma j}) \omega^\sigma \wedge \omega^j \Rightarrow$ if $\Gamma^k_{\sigma j} = \Gamma^k_{j\sigma}$

$$\sum_{j<\sigma} 2R^\nu_{\mu j\sigma} \omega^\sigma \wedge \omega^j = \sum_{j<\sigma} (\partial_j \Gamma^\nu_{\mu\sigma} - \partial_\sigma \Gamma^\nu_{\mu j} + \Gamma^\nu_{\mu\sigma} \Gamma^k_{kj} - \Gamma^\nu_{\mu j} \Gamma^k_{k\sigma}) \omega^\sigma \wedge \omega^j$$

(D) If this is the case say ∇ is symmetric

OR: use (*)

so

$\rightarrow$ ||

$2R^\nu_{\mu j\sigma}$

(There's no "2" in books. Not clear why.)

(P) If g is a metric $\exists!$ symmetric connection ∇ s.t. $\nabla g = 0$.
(call ∇ the Levi-Civita conn.)

Pf Uniqueness. $\nabla g = 0 \Rightarrow 0 = \nabla_\sigma g_{\mu\nu} \stackrel{(*)}{=} \partial_\sigma g_{\mu\nu} - \Gamma^\rho_{\mu\sigma} g_{\rho\nu} - \Gamma^\rho_{\nu\sigma} g_{\mu\rho} \Rightarrow \partial_\sigma g_{\mu\nu} = \Gamma^\rho_{\mu\sigma} g_{\rho\nu} + \Gamma^\rho_{\nu\sigma} g_{\mu\rho}$

$$\Rightarrow \partial_\nu g_{\sigma\mu} + \partial_\mu g_{\sigma\nu} - \partial_\sigma g_{\mu\nu} = \cancel{\Gamma^\rho_{\nu\sigma} g_{\mu\rho}} + \Gamma^\rho_{\nu\mu} g_{\sigma\rho} + \cancel{\Gamma^\rho_{\mu\sigma} g_{\rho\nu}} + \Gamma^\rho_{\mu\nu} g_{\sigma\rho} - \cancel{\Gamma^\rho_{\mu\sigma} g_{\rho\nu}} - \cancel{\Gamma^\rho_{\nu\sigma} g_{\mu\rho}}$$

$$= 2 g_{\sigma i} \Gamma^i_{\mu\nu}$$

$$\Rightarrow \frac{1}{2} g^{\tau\sigma} \{ \partial_\nu g_{\sigma\mu} + \partial_\mu g_{\sigma\nu} - \partial_\sigma g_{\mu\nu} \} = g^{\tau\sigma} g_{\sigma i} \Gamma^i_{\mu\nu} = \delta^\tau_i \Gamma^i_{\mu\nu} = \Gamma^\tau_{\mu\nu}$$

if g^{ij} inv of g_{ij}

Existence: define $\Gamma^\tau_{\mu\nu}$ by & check

F^α_μ

(*) So the $n \times n$ matrix of 2-forms $(R^\alpha_{\mu\rho\sigma} \omega^\rho \wedge \omega^\sigma)_{\alpha\mu}$ is the curvature of ∇

(check if this is the right contraction !!)

① Ricci tensor $R_{\mu\nu} = R^{\lambda}_{\mu\lambda\nu}$; Ricci curvature $R = g^{\mu\nu} R_{\mu\nu} \in \mathcal{O}(M)$
 $R_{\mu\nu} \omega^{\mu} \otimes \omega^{\nu}$

bicanonical form $\eta_g = (\det g_{\mu\nu}) (\omega^1 \wedge \omega^2 \wedge \dots \wedge \omega^d) \otimes^2 \in (\Omega^d)^{\otimes 2} = K^{\otimes 2}$, $d = \dim M$
 (correctly def b/c if $g_{\mu\nu} \omega^{\mu} \otimes \omega^{\nu} = \tilde{g}_{\mu\nu} \tilde{\omega}^{\mu} \otimes \tilde{\omega}^{\nu}$, $\tilde{\omega}^{\mu} = \gamma^{\mu}_{\tilde{\mu}} \omega^{\tilde{\mu}}$
 $\Rightarrow \text{RHS} = \tilde{g}_{\mu\nu} \gamma^{\mu}_{\tilde{\mu}} \gamma^{\nu}_{\tilde{\nu}} \omega^{\tilde{\mu}} \otimes \omega^{\tilde{\nu}} \Rightarrow g = \gamma \tilde{g} \gamma^T \Rightarrow \det g = (\det \tilde{g}) (\det \gamma)^2$
 $\& \tilde{\omega}^1 \wedge \dots \wedge \tilde{\omega}^d = \det \gamma \cdot \omega^1 \wedge \dots \wedge \omega^d$

a canonical form is a form $\omega \in \Omega^d K$ s.t. $\omega \otimes^2 = \pm \eta_g$ (it may not $\exists$ in alg setting; it $\exists$ on double cover; it always $\exists$ in C^∞ case & is "can")
 b/c $\det g_{\mu\nu}$ has const sign & $|\det g_{\mu\nu}|^{1/2}$ is C^∞)

Einstein-Hilbert Lagr. $L_{HE} = R_{HE} \in \Omega^d(M)$

Hodge star (in C^∞ case) $* = \frac{1}{g}$

$* \left(\frac{1}{p!} \varphi_{i_1 \dots i_p} \omega^{i_1} \wedge \dots \wedge \omega^{i_p} \right) = \frac{1}{(d-p)! p!} \varepsilon_{i_1 \dots i_p} |\det g_{\mu\nu}|^{1/2} \varphi_{j_1 \dots j_p} g^{j_1 i_1} \dots g^{j_p i_p} \omega^{j_1} \wedge \dots \wedge \omega^{j_p}$

(Note Yang-Mills for $* = \frac{1}{g}$ becomes $L_{YM}(\Xi, A, g)$ while $L_{HE} = L_{HE}(g)$ (check again!))
Gauge intpr of Einst. theory

