

On isospectral deformations of an inhomogeneous string

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Abstract

In this paper we consider a class of isospectral deformations of the inhomogeneous string boundary value problem. The deformations considered are generalizations of the isospectral deformation that has arisen in connection with the Camassa-Holm equation for the shallow water waves. It is proved that these new isospectral deformations result in evolution equations on the mass density whose form depends on how the string is tied at the endpoints. Moreover, it is shown that the evolution equations in this class linearize on the spectral side and hence can be solved by the inverse spectral method. In particular, the problem involving a mass density given by a discrete finite measure and arbitrary boundary conditions is shown to be solvable by Stieltjes' continued fractions.

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1 Introduction

The wave equation $\frac{1}{c^2} f_{\tau\tau} - \nabla^2 f = 0$ is a prototype of hyperbolic problems well studied since the time of Euler. One of its many applications, well known to any science undergraduate student, is the one dimensional version of this equation $\frac{1}{c^2} f_{\tau\tau} - f_{xx} = 0$, $0 < x < L$, describing the motion of a vibrating string of length L with the amplitude $f(x, \tau)$ at the point x and time τ . Usually, the equation is also equipped with boundary conditions expressing the way the endpoints of the string are behaving; the most common types are, of course, the case of the string with both ends fixed (i.e. $f(0, \tau) = f(L, \tau) = 0$) or the case of the string with sliding endpoints (i.e. $f_x(0, \tau) = f_x(L, \tau) = 0$). The two cases are sometimes referred to as the Dirichlet and the Neumann strings respectively. The coefficient c^2 has the physical dimension of the velocity squared but more appropriately it depends on the mass density of the medium. In general, in the case of an inhomogeneous medium, this coefficient is position dependent. The string counterpart of this is an *inhomogeneous string* equation: $\frac{1}{c^2(x)} f_{\tau\tau} - f_{xx} = 0$, $0 < x < L$ with the position dependent speed $c^2(x) = \frac{T}{\rho(x)}$, where T is the tension and ρ the lineal mass density of the string, which leads, after separation of variables $f(x, \tau) = \cos(\sqrt{z}\tau) v(x)$, to

$$-v_{xx} = \frac{z}{T} \rho(x) v. \quad (1.1)$$

We will choose the tension $T = 1$ and for simplicity of presentation we will restrict our attention to the string problem with $L = 1$. The general string boundary value problem then reads:

$$-v_{xx} = z\rho(x)v, \quad 0 < x < 1, \quad v_x(0) - hv(0) = 0, \quad v_x(1) + Hv(1) = 0, \quad (1.2)$$

where, traditionally, $h = 0$ corresponds to the Neumann boundary condition on the left end of the string, while $h = \infty$ corresponds to the Dirichlet condition on the left end, and the same convention applies to H and the right end of the string. We will subsequently refer to equation (1.2) as the string equation, although the name of the Helmholtz equation is also associated with this very equation.

The mathematical pedigree of equation (1.2) parallels its physical significance: (1.2) impacted the development of the Fourier methods, continued fractions, and subsequently led to the fundamental progress associated with distributional (weak) solutions to partial differential equations. For this paper, however, the most relevant is M.G. Krein's pioneering work on the inverse problem for the inhomogeneous string [11, 18, 19]. The inverse problem for the inhomogeneous string, in its simplest formulation, amounts to the reconstruction of the mass density ρ from the spectrum of the boundary value problem. It has been known since the time of Borg [8] that the mass density of the string is not uniquely determined by the spectrum alone. It was M.G. Krein who shifted the emphasis to the reconstruction of the string density from the spectral function, sometimes also called the Weyl function. In the simplest case of the Dirichlet spectrum the corresponding Weyl function is $W(z) = \frac{v_x(1; z)}{v(1; z)}$, where $v(x; z)$ satisfies the initial value problem: $-v_{xx} = z\rho(x)v$, $0 < x < 1$, $v(0; z) = 0$, $v_x(0; z) = 1$. The zeros of

the denominator of the Weyl function correspond to the spectrum of the Dirichlet boundary value problem, while the zeros of the numerator correspond to the spectrum of the boundary value problem with the Dirichlet condition on the left and the Neumann condition on the right. The inverse problem for the string appears in many different areas of science and engineering such as the geophysical inverse problems [2] or magnetotelluric inversion [15, 26], to mention just a few. The reader might want to consult a very accessible account of inverse problems relevant to engineering in the book by G.M.L. Gladwell [13].

Since in general the knowledge of a single spectrum does not determine the string density, a natural question is to determine *isospectral* deformations of the string, i.e. changes in the density that leave invariant the spectrum of a given boundary value problem.

Some results dealing with isospectral strings are presented in [13] and also in the works of H.P.W. Gottlieb [14] and in the references therein. Our interest lies in characterizing isospectral strings in terms of *nonlinear evolution equations* on the mass density ρ . The first contribution to this line of research, as far as we know, goes back to the series of interesting papers by P.C. Sabatier [21–23] (see Example 2.2). Our approach, however, is more influenced by the later work on the Camassa-Holm equation (CH) [9]:

$$m_t = m_x u + 2u_x m, \quad m = u - u_{xx}.$$

It was discovered in the late nineties in [4] that the CH equation can be viewed as an isospectral deformation of the inhomogeneous string with Dirichlet boundary conditions. This led, among other things, to the construction of a certain class of explicit solutions (peakons) in [5] by adapting an elegant approach of Stieltjes [24]. Some of these developments are reviewed in [7].

The present paper addresses the following questions:

1. consider a general class of deformations of the boundary value problem (1.2) postulated to be described by $v_t = av + bv_x$, where all functions depend on x , the deformation parameter t , and the spectral parameter z appearing in equation (1.2). Can a and b be suitably chosen for the boundary value problem (1.2) to remain isospectral?
2. which of the isospectral nonlinear evolution equations on the mass density give rise to a linear evolution of the appropriate Weyl function?

In the body of the paper we give a definite, affirmative, answer to the first question under the assumption that b is a rational function of z , regular at ∞ and possessing only simple poles which, for technical reasons, we assume to be located on the positive part of the real axis.

Even though the second question is less natural from the point of view of real life applications it is nevertheless of considerable mathematical interest because the linear evolution of the Weyl function leads, via the spectral/inverse scattering approach, to exact formulas which can be of great benefit, for example, in testing numerical methods. Furthermore, by constructing explicit solutions to nonlinear partial differential equations, one frequently gets valuable insight into the expected behaviour of a larger class of solutions after applying, for example, density arguments if such are available. The research presented in this paper provides a partial answer to the question of linearization for the the following reason: we rely on a particular scheme of linearization originally known from the inverse scattering approach to the KdV equation and subsequently applied to a multitude of integrable partial differentiable equations. This method hinges in an essential way on the linearization of dynamics

in the asymptotic region in space, eventually leading to a linear evolution of the scattering data. In summary, we answer the question of linearization only within the paradigm of the inverse scattering approach.

The paper is organized as follows: in section 2 we analyze what we call the Zakharov-Shabat deformations (ZS) modeled after those occurring in the theory of integrable systems (see, for example, [12]); section 3 contains a preliminary analysis of these flows, we derive a criterion for the preservation of boundary conditions under the deformations and provide evidence that the ZS isospectral deformations of the string arise from the class of flows analytic around $z = \infty$; section 4 contains, in Proposition 4.1, the description of these isospectral deformations in terms of evolution equations on the spectral data; in section 5 we prove that the deformation, rational in the spectral variable, regular at infinity, with simple poles at positive $\epsilon_1, \epsilon_2, \dots, \epsilon_M$, can be chosen to render equation (1.2) not only isospectral but also such that the evolution of the spectral data is linear; in section 6 we show explicitly how the deformations work for discrete mass densities; finally, in Appendix 8 we provide details behind the modifications of the formalism when the mass densities are not smooth but rather finite discrete measures, while in Appendix 9 we explain the main steps illustrating the relation of the string deformation equations with the CH equation.

2 Inhomogeneous string and its deformations

We will make the general assumption that ρ is a positive measure and $h, H \geq 0$ to ensure that the spectrum of the boundary value problem (1.2) is positive. However, to simplify the technical aspects of the paper, we will focus on two extreme cases:

1. ρ is in $C^1([0, 1])$,
2. ρ is a finite, discrete measure, i.e. $\rho = \sum_{j=1}^N m_j \delta_{x_j}$ where δ_{x_j} is the Dirac measure centered at $0 < x_j < 1$,

leaving other scenarios of the smoothness of the mass density for future work.

We begin by writing the string equation as a first order system:

$$V_x = \begin{bmatrix} 0 & 1 \\ -z\rho & 0 \end{bmatrix} V, \quad \text{where } V = \begin{bmatrix} v \\ v_x \end{bmatrix}, \quad (2.1)$$

and define deformations of the string equation as variations of the mass density ρ expressed as a function of an external parameter t which, we emphasize, is not a physical time, even though we will occasionally refer to it as time. Our goal is to describe a class of deformations which leave invariant the spectrum of the string boundary value problem (1.2); we will refer to these deformations as *isospectral deformations* and will search for them among deformations of the Zakharov-Shabat (ZS) type.

Definition 2.1 (ZS deformations).

V in equation (2.1) is postulated to depend on the deformation parameter t subject to

$$V_t = \begin{bmatrix} a & b \\ c & d \end{bmatrix} V \quad (2.2)$$

where a, b, c, d are functions of x, t, z , chosen so as to ensure the compatibility of equations (2.1) and (2.2).

When ρ, a, b, c, d are sufficiently smooth in x and t (for example all in C^1) the equations (2.1) and (2.2) are compatible provided the *Zero Curvature condition* :

$$\begin{bmatrix} 0 & 1 \\ -z\rho & 0 \end{bmatrix}_t - \begin{bmatrix} a & b \\ c & d \end{bmatrix}_x + \left[\begin{bmatrix} 0 & 1 \\ -z\rho & 0 \end{bmatrix}, \begin{bmatrix} a & b \\ c & d \end{bmatrix} \right] = 0, \quad (2.3)$$

holds, which, upon closer inspection, implies that one can parametrize a, c, d by the function b . The result of a simple computation is

$$a = -\frac{1}{2}b_x + \beta, \quad c = -\frac{1}{2}b_{xx} - z\rho b, \quad d = \frac{1}{2}b_x + \beta, \quad (2.4)$$

where β is a constant (in x), whose choice, to be discussed later, is dictated by the boundary conditions. Then we substitute these into the equation containing ρ_t to get the following *evolution deformation equation*

$$z\rho_t = \frac{1}{2}b_{xxx} + z\rho_x b + 2z\rho b_x, \quad (2.5)$$

which we will write as

$$z\rho_t = \frac{1}{2}b_{xxx} + z\mathcal{L}_\rho b$$

with

$$\mathcal{L}_\rho = \rho \frac{\partial}{\partial x} + \frac{\partial}{\partial x} \rho \quad (2.6)$$

denoting the second and third term in (2.5). In general, one cannot say much more about equation (2.5) unless some additional structure in the spectral variable z is brought to bear. For example, assuming the functions b, b_x , etc., to have Laurent expansions in z , that is to say $b = \sum b_n z^n$, $b_x = \sum b_{n,x} z^n$, and $b_{xxx} = \sum b_{n,xxx} z^n$, then collecting like powers of z in (2.5) gives

$$\rho_t = \frac{1}{2}b_{1,xxx} + \mathcal{L}_\rho b_0, \quad \frac{1}{2}b_{n,xxx} + \mathcal{L}_\rho b_{n-1} = 0, \quad \text{for } n \neq 1. \quad (2.7)$$

If one assumes b to be a rational function

$$b(z) = b_0 + \sum_{k=1}^M \frac{b_{-1}^{(k)}}{z + \epsilon_k}, \quad \epsilon_k \neq \epsilon_j, \quad k \neq j \text{ and } \epsilon_k > 0 \quad (2.8)$$

then the deformation equation (2.5) reads:

$$\rho_t = \mathcal{L}_\rho b_0, \quad (2.9a)$$

$$b_{0,xxx} + \sum_{k=1}^M \frac{b_{-1,xxx}^{(k)}}{\epsilon_k} = 0, \quad \frac{1}{2}b_{-1,xxx}^{(k)} - \epsilon_k \mathcal{L}_\rho b_{-1}^{(k)} = 0, \quad k = 1 \dots M. \quad (2.9b)$$

Let us consider a few specific examples of deformation equations.

Example 2.2. Taking $b = zb_1 + b_0$ one obtains

$$\rho_t = \frac{1}{2}b_{1,xxx} + \mathcal{L}_\rho b_0, \quad \mathcal{L}_\rho b_1 = 0, \quad b_{0,xxx} = 0. \quad (2.10)$$

This deformation was mentioned already in [22] where this evolution equation was identified as the *Harry-Dym equation* (HD).

Example 2.3. Similarly, the case $b = b_0 + \frac{b_{-1}}{z}$ results in

$$\rho_t = \mathcal{L}_\rho b_0, \quad \frac{1}{2}b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad b_{-1,xxx} = 0. \quad (2.11)$$

The relation of this deformation to the string with Dirichlet boundary conditions was established in [5] and, in this special case, the evolution equation (2.11) (after a Liouville transformation) corresponds to the Camassa-Holm (CH) equation [9] mentioned earlier.

Finally we have the rational model, the main focus of this paper.

Example 2.4. We put $M = 1$, set $\epsilon_1 = \epsilon$, $b_{-1}^{(1)} = b_{-1}$ in (2.8) and rearrange equations to obtain

$$\rho_t = \mathcal{L}_\rho b_0, \quad \frac{1}{2}b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad \epsilon b_{0,xxx} + b_{-1,xxx} = 0. \quad (2.12)$$

Remark 2.5. The idea behind the use of the parameter ϵ is that we want to have control over the singularity of b , in particular we do not want to have a singularity at a point of the spectrum of the string. Moreover, we plan to compute the limit $\epsilon \rightarrow 0$ in the rational model. The presentation of example 2.4 suggests that, at least formally, the rational model goes over to the one discussed in example 2.3. Finally, the relevance of the rational model, in addition to the fact that it goes in the limit $\epsilon \rightarrow 0$ to other known cases, is that it admits a natural resolution of the differential constraints in terms of products of eigenfunctions, exemplified by Theorems 5.8, 5.13, or in the general case of rational flows (2.8), by Theorem 5.9.

In the remainder of this section we discuss one type of non-smooth mass densities, exemplified by $\rho = \sum_{j=1}^N m_j \delta_{x_j}$.

Away from the support of ρ we are dealing with the smooth case so our previous discussion applies. However, on the support of ρ we need to apply distributional (in the sense of theory of distributions) calculus. In Appendix 8, we analyze the deformation equation on the support of ρ and show that the deformation equation will take the form of the distributional equation

$$zD_t \rho = \frac{1}{2}D_x^3 b + z\mathcal{L}_\rho b \quad (2.13)$$

where D_t, D_x are distributional derivatives in t, x respectively while $\mathcal{L}_\rho f = D_x(\rho f) + \rho \langle f_x \rangle$ for any continuous piecewise smooth function f with $\langle f_x \rangle$ denoting the average function (the arithmetic average of the right-hand and left-hand limits).

Example 2.6. For the rational flow $b = b_0 + \frac{b_{-1}}{z+\epsilon}$ and $\rho = \sum_{j=1}^N m_j \delta_{x_j}$ the deformation equations restricted to the support of ρ read:

$$\begin{aligned} D_t \rho &= \sum_{j=1}^N \dot{m}_j \delta_{x_j} - m_j \dot{x}_j \delta'_{x_j}, \\ z m_j \dot{x}_j &= -\frac{1}{2} [b_x](x_j) - z m_j b(x_j), \\ z \dot{m}_j &= \frac{1}{2} [b_{xx}](x_j) + z m_j \langle b_x \rangle(x_j), \end{aligned} \quad (2.14)$$

where $[f](x_j)$ is the jump of f at x_j . Writing these equations in terms of components b_0 and b_{-1} we obtain:

$$\dot{x}_j = -b_0(x_j), \quad -\epsilon [b_{0,x}](x_j) = [b_{-1,x}](x_j), \quad \frac{1}{2} [b_{-1,x}](x_j) = \epsilon m_j b_{-1}(x_j) \quad (2.15)$$

$$\dot{m}_j = m_j \langle b_{0,x} \rangle(x_j), \quad -\epsilon [b_{0,xx}](x_j) = [b_{-1,xx}](x_j), \quad \frac{1}{2} [b_{-1,xx}](x_j) = \epsilon m_j \langle b_{-1,x} \rangle(x_j), \quad (2.16)$$

producing, in the formal limit $\epsilon \rightarrow 0$ (corresponding to $b = b_0 + \frac{b_{-1}}{z}$), the dynamical system similar to the one known from the CH theory:

$$\dot{x}_j = -b_0(x_j), \quad \dot{m}_j = m_j \langle b_{0,x} \rangle(x_j). \quad (2.17)$$

These are only preliminary computations which will be fully justified once the existence of the limit for b_{-1} and b_0 is proved.

3 Boundary conditions

The boundary conditions (1.2) can be written as

$$[-h \quad 1] V(x=0) = 0, \quad [H \quad 1] V(x=1) = 0.$$

We require that these boundary conditions hold during the deformation of the string, in other words we require

$$[-h \quad 1] V_t(x=0) = 0, \quad [H \quad 1] V_t(x=1) = 0,$$

or

$$\begin{aligned} [-h \quad 1] \begin{bmatrix} a & b \\ c & d \end{bmatrix} V(x=0) &= 0, \\ [H \quad 1] \begin{bmatrix} a & b \\ c & d \end{bmatrix} V(x=1) &= 0, \end{aligned} \quad (3.1)$$

where a , c , and d have been determined in (2.4) in terms of the function b , thus leading to the explicit, b -dependent, form of (3.1).

Lemma 3.1. *Suppose v satisfies the string boundary value problem (1.2) at time t . Then the ZS deformation leaves the boundary conditions invariant provided*

$$\begin{aligned} \frac{1}{2} b_{xx}(0) - h b_x(0) + (h^2 + z\rho(0))b(0) &= 0, \\ \frac{1}{2} b_{xx}(1) + H b_x(1) + (H^2 + z\rho(1))b(1) &= 0. \end{aligned} \quad (3.2)$$

In particular, if there is no mass at the endpoints then

$$\begin{aligned}\frac{1}{2}b_{xx}(0) - hb_x(0) + h^2b(0) &= 0, \\ \frac{1}{2}b_{xx}(1) + Hb_x(1) + H^2b(1) &= 0,\end{aligned}\tag{3.3}$$

with the proviso that $b(0) = 0$ if $h = \infty$, likewise $b(1) = 0$ if $H = \infty$.

Remark 3.2. We emphasize that condition (3.2) has to hold not only in t but also in z , resulting in conditions on the components of b . The latter leads to severe restrictions on admissible choices of b , for example, eliminating the polynomial flows as the example below illustrates.

We briefly analyze a few special cases of Lemma 3.1.

Example 3.3 (Neumann and Dirichlet Conditions). The Neumann-Neumann boundary conditions have $h = H = 0$, hence $b_{xx}(0) = b_{xx}(1) = 0$, in the case of no masses at the end points. In the Dirichlet-Dirichlet case, regardless of the presence of masses at the endpoints, $h = H = \infty$, hence $b(0) = b(1) = 0$.

We illustrate the relevance of Lemma 3.1 in highlighting the difference brought about by specific analytic properties of b at $z = \infty$. In particular, we would like to gauge to which extent it is possible to impose the boundary conditions (3.2) on b without violating the differential equation (2.5). To appreciate the task we will analyze two cases.

Polynomial flows: $\mathbf{b} = \mathbf{b}_1\mathbf{z} + \mathbf{b}_0$

The differential equations (2.10) yield solutions of the form

$$b_0(x) = Ax^2 + Bx + C, \quad b_1(x) = \frac{D}{\rho(x, t)^{1/2}}.\tag{3.4}$$

Suppose we impose Dirichlet-Dirichlet boundary conditions. Then Lemma 3.1 implies $b_0(0) = b_0(1) = 0$ and $b_1(0) = b_1(1) = 0$, hence $b_0(x) = Ax(1 - x)$ while by the second condition either $D = 0$ or

$$b_1(0) = b_1(1) = \frac{D}{\rho(0, t)^{1/2}} = \frac{D}{\rho(1, t)^{1/2}} = 0.$$

Thus the polynomial flow does not preserve Dirichlet-Dirichlet boundary conditions unless the density is singular at the endpoints. Consequently, for general densities, $D = 0$, hence $b_1(x) = 0$, and the only polynomial flow preserving Dirichlet-Dirichlet boundary conditions will be linear: $\rho_t = \mathcal{L}_\rho b_0$, $b_0 = Ax(1 - x)$. To deal with boundary conditions other than Dirichlet-Dirichlet boundary conditions we observe that by inspection equation (3.2) implies $\frac{1}{2}b_{i,xx}(0) - hb_{i,x}(0) + h^2b_i(0) + \rho(0)b_{i-1}(0) = 0$. So, in the case of $i = 2$, if there is a mass at the left endpoint, $b_1(0) = 0$ which implies $D = 0$ unless the density is singular at $x = 0$. Thus in the regular case again $b_1(x) = 0$ and we conclude that for the general boundary conditions and general mass densities the linear flow above is the only flow preserving these boundary conditions. This argument can easily be generalized to the case of polynomial flows $\sum_{n=0}^N b_n z^n$. We conclude this example by pointing out that because of the form of b_1 which involves both taking the root and inverting of ρ , one does not expect the situation to improve for ρ given by discrete measures.

Next, we illustrate how the behaviour of b at $z = \infty$ impacts analysis.

Flows polynomial in $1/z$: $\mathbf{b} = \mathbf{b}_0 + \frac{\mathbf{b}_{-1}}{z}$

For simplicity we will only analyze the case of vanishing mass density at the endpoints. Recall the system of differential equations for the case $b = b_0 + \frac{b_{-1}}{z}$ (equation (2.11)):

$$\rho_t = \mathcal{L}_\rho b_0, \quad \frac{1}{2}b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad \frac{1}{2}b_{-1,xxx} = 0.$$

The last equation gives $b_{-1}(x) = \frac{1}{2}\alpha_1 x^2 + \alpha_2 x + \alpha_3$. The boundary conditions on b_{-1} given by equations (3.3) imply:

$$\begin{bmatrix} \frac{1}{2} & -h & h^2 \\ \frac{1}{2} + H + \frac{1}{2}H^2 & H + H^2 & H^2 \end{bmatrix} \begin{bmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}. \quad (3.5)$$

These equations, by virtue of representing two planes going through the origin, will always have a solution.

Now, as to b_0 we observe that one can write b_0 in terms of ρ and b_{-1} by integrating

$$\begin{aligned} b_{0,xx} &= -2 \int \rho_x b_{-1} dx - 4 \int \rho b_{-1,x} dx + C_1 = -2\rho b_{-1} - 2 \int \rho b_{-1,x} dx + C_1, \\ b_{0,x} &= -2 \int (\rho b_{-1} + \int \rho b_{-1,x} dx) dx + C_1 x + C_2 := f(x) + C_1 x + C_2, \\ b_0 &= \int f(x) dx + \frac{1}{2}C_1 x^2 + C_2 x + C_3. \end{aligned} \quad (3.6)$$

Observe that we can always choose $f(x)$ above in such a way that $b_0(0) = C_3$, $b_{0,x}(0) = C_2$, $b_{0,xx}(0) = C_1$ and then set $b_0(x) = \int_0^x f(\xi) d\xi + \frac{1}{2}C_1 x^2 + C_2 x + C_3$. Then the boundary conditions on b_0 given by equations (3.3) imply:

$$\begin{bmatrix} \frac{1}{2} & -h & h^2 \\ \frac{1}{2} + H + \frac{1}{2}H^2 & H + H^2 & H^2 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \\ C_3 \end{bmatrix} = \begin{bmatrix} 0 \\ C \end{bmatrix}, \quad (3.7)$$

where $C = \int_0^1 \rho(x) b_{-1,x}(x) dx - Hf(1) - H^2 \int_0^1 f(x) dx$. This system of equations describes two planes, one of which goes through the origin. These planes will always intersect unless they are parallel, in which case the normal vectors are proportional. For this to happen

$$\begin{bmatrix} \frac{1}{2} & -h & h^2 \end{bmatrix} = k \begin{bmatrix} \frac{1}{2} + H + \frac{1}{2}H^2 & H + H^2 & H^2 \end{bmatrix},$$

for some constant k . Since $0 \leq h, H$ we see that the only case when this occurs is the case $h = H = 0$. However, if $h = H = 0$ then equation (3.5) implies that b_{-1} is a linear function in x . Moreover, in this case, for equation (3.7) to have a solution, C must vanish. Thus $\int_0^1 \rho(x) b_{-1,x} dx = 0$ and it suffices to choose b_{-1} to be a constant to satisfy (3.7). The remaining cases involving either $h = \infty$ or $H = \infty$ or $h = H = \infty$ can be analyzed in the same way. In summary, for all $0 \leq h, H \leq \infty$ one can satisfy conditions (3.3). Finally, in the case of Neumann-Neumann conditions ($h = H = 0$) one recovers the Hunter-Saxton equation [6, 16]. We will continue this analysis in later sections, in particular we will give a uniform construction of b valid for all boundary conditions (see theorem 5.19).

4 Isospectrality for general boundary conditions

In this section we check directly that the b flows, subject to (3.3), are isospectral for the boundary value problem (1.2). To simplify analysis, and also guided by the inverse scattering technique, we will make the following assumption.

Assumption

From this point onwards we will be assuming that there is no mass in some open neighborhoods I_0, I_1 of the left endpoint, right endpoint, respectively.

Our starting point is again (2.5). We assume in this section that conditions (3.2) hold to ensure that the ZS deformations leave the boundary conditions invariant by Lemma 3.1. The assumption on the absence of masses around the endpoints is in force.

On I_0

$$v(x) = \begin{cases} hx + 1 & \text{if } 0 \leq h < \infty, \\ x + \frac{1}{h} & \text{if } 0 < h \leq \infty, \end{cases} \quad (4.1)$$

satisfies equation (1.2). Moreover, since there is no mass on I_0 , $b_{xxx} = 0$ by equation (2.5). Hence b is quadratic in x on I_0 . Evaluating the deformation $v_t = (\beta - \frac{1}{2}b_x)v + bv_x$ for v satisfying equations (4.1) and (3.2) yields:

$$\beta = \begin{cases} \frac{1}{2}b_x(0) - hb(0) & \text{if } 0 \leq h < \infty, \\ -\frac{1}{2}b_x(0) + \frac{1}{2h}b_{xx}(0) & \text{if } 0 < h \leq \infty. \end{cases} \quad (4.2)$$

Now we turn to analyzing the behaviour on I_1 , finding from (1.2) that

$$v(x) = A_1(t, z)(x-1) + B_1(t, z), \quad (4.3)$$

implying that the spectrum, determined by the roots of $v_x(1) + Hv(1) = 0$ will be defined by the roots of $A_1(t, z) + HB_1(t, z) = 0$ or $\frac{1}{H}A_1(t, z) + B_1(t, z) = 0$ according to whether $0 \leq H < \infty$ or $0 < H \leq \infty$ respectively. As on I_0 we similarly find that on I_1

$$b = \frac{1}{2}b_{xx}(1)(x-1)^2 + b_x(1)(x-1) + b(1);$$

suppressing the dependence of coefficients on t and z to avoid a cluttered notation. Collecting like powers of $(x-1)$ in the deformation $v_t = (\beta - \frac{1}{2}b_x)v + bv_x$ as well as using equation (3.3) gives

$$\dot{A}_1 = (\beta + \frac{1}{2}b_x(1))A_1 - \frac{1}{2}b_{xx}(1)B_1 = (\beta + \frac{1}{2}b_x(1))A_1 + H(b_x(1) + Hb(1))B_1, \quad (4.4)$$

$$\dot{B}_1 = (\beta - \frac{1}{2}b_x(1))B_1 + b(1)A_1 = (\beta - \frac{1}{2}b_x(1))B_1 - \frac{1}{H^2}(Hb_x(1) + \frac{1}{2}b_{xx}(1))A_1. \quad (4.5)$$

Hence when $0 \leq H < \infty$

$$\dot{A}_1 + H\dot{B}_1 = \left(\beta + \frac{1}{2}b_x(1) + Hb(1)\right)(A_1 + HB_1),$$

or when $0 < H \leq \infty$

$$\frac{1}{H}\dot{A}_1 + \dot{B}_1 = \left(\beta - \frac{1}{2}b_x(1) - \frac{1}{2H}b_{xx}(1)\right)\left(\frac{1}{H}A_1 + B_1\right)$$

holds. Thus under the deformation, $v_x(1) + Hv(1) = 0$ will persist, hence the ZS deformation is indeed isospectral. However, the extent to which the evolution of the spectral data encoded in (4.4), (4.5) is simpler than the original evolution equation on ρ is far from certain, for example the coefficient β will in general depend (nonlocally) on ρ and, in principle, could even depend on the deformation time t . The remainder of the paper is devoted to the goal of linearizing the evolution on the spectral side. Towards this end we will rephrase equations (4.4) and (4.5) by differently grouping terms on the right hand side of respective equations.

Proposition 4.1. *Let v satisfy the string boundary value problem (1.2) for the mass density ρ with mass gaps around the endpoints, $0 \leq h, H \leq \infty$ and let the ZS deformation satisfy the deformation equation (2.5) with b subject to (3.3). Let us define a function K by*

$$K\beta = \begin{cases} \frac{1}{2}b_x(1) + Hb(1) & \text{if } 0 \leq H < \infty, \\ -\frac{1}{2}b_x(1) - \frac{1}{2H}b_{xx}(1) & \text{if } 0 < H \leq \infty, \end{cases} \quad (4.6)$$

where β is given by (4.2). Then the deformation yields the following evolution equations of the spectral data A_1, B_1 :

1. For $0 < H < \infty$:

$$\dot{A}_1 + H\dot{B}_1 = (K + 1)\beta[A_1 + HB_1], \quad (4.7a)$$

$$\dot{A}_1 - H\dot{B}_1 = (1 - K)\beta[A_1 - HB_1] + b_x(1)[A_1 + HB_1], \quad (4.7b)$$

2. For $H = 0$:

$$\dot{A}_1 = (1 + K)\beta A_1, \quad (4.8a)$$

$$\dot{B}_1 = (1 - K)\beta B_1 + b(1)A_1, \quad (4.8b)$$

3. For $H = \infty$:

$$\dot{A}_1 = (1 - K)\beta A_1 - \frac{1}{2}b_{xx}(1)B_1, \quad (4.9a)$$

$$\dot{B}_1 = (1 + K)\beta B_1. \quad (4.9b)$$

Remark 4.2. The objective of the next section is to show that one can choose the field b , without violating (3.3), in such a way that $K = -1$ and β is invariant under t , thereby making it computable from $\rho(t = 0)$.

5 Isospectrality and linearization for rational flows and general boundary conditions

As we observed in section 3 polynomial ZS flows are not suitable deformations of a string boundary value problem for mass densities regular at the endpoints, and also for those for which the mass density has zeros (see (3.4)). On the other hand, the flow of the CH type, namely $b = b_0 + \frac{b-1}{z}$, works fine in

either case. So the most natural first step towards a generalization is simply to try $b = b_0 + \sum_{j=1}^n \frac{b_{-j}}{z^j}$. As stated earlier, (see (2.7)) one then obtains

$$\rho_t = \mathcal{L}_\rho b_0, \quad \frac{1}{2}b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad \frac{1}{2}b_{-n,xxx} + \mathcal{L}_\rho b_{-(n+1)} = 0, \quad 1 \leq n.$$

The truncation at $n = 2$, which amounts to setting $b_{-n} = 0$ starting with $n = 2$, produces the CH type system, with two interacting components b_0 and b_{-1} . There is however another scenario in which one solves the constraints $\frac{1}{2}b_{-n,xxx} + \mathcal{L}_\rho b_{-(n+1)} = 0$, $1 \leq n$, not by truncation but by a sort of "quasi-periodicity", positing $b_{-(n+1)} = w b_{-n}$, $2 \leq n$, for some proportionality constant w (taken to be $-\epsilon$ in the remainder of this paper). Then out of the infinite set of equations one needs only two, namely,

$$\frac{1}{2}b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad \frac{1}{2}b_{-1,xxx} + w \mathcal{L}_\rho b_{-1} = 0,$$

and the theory remains described by two fields only. This is one of the motivations for studying the rational model.

We will argue now that the rational flows are free of the deficiencies present for polynomials flows described above; in particular the vanishing of ρ is no longer an issue. To this end we will analyze the simplest rational flow with one pole. Thus $b = b_0 + \frac{b_{-1}}{z+\epsilon}$, $\epsilon > 0$ (equation (2.8)), For convenience we recall the differential equations (2.12) in the special case of one pole, writing them in a form convenient for further analysis:

$$\rho_t = \mathcal{L}_\rho b_0, \tag{5.1a}$$

$$b_{0,xxx} + \frac{b_{-1,xxx}}{\epsilon} = 0, \quad \frac{1}{2}b_{-1,xxx} - \epsilon \mathcal{L}_\rho b_{-1} = 0. \tag{5.1b}$$

We begin our analysis by formulating three auxiliary lemma, the first of which can be found for example in [1], and is discussed in [3], [25].

Lemma 5.1. *Let $\lambda \in \mathbb{C}$, $\rho \in C^1([0, 1])$ and suppose ϕ and ψ are solutions of $f_{xx} - \lambda \rho f = 0$. Then ϕ^2 , $\phi\psi$, and ψ^2 are solutions of $\frac{1}{2}g_{xxx} - \lambda \mathcal{L}_\rho g = 0$. If, in addition, ϕ and ψ are linearly independent, then ϕ^2 , $\phi\psi$, and ψ^2 span the solution space of $\frac{1}{2}g_{xxx} - \lambda \mathcal{L}_\rho g = 0$.*

Proof. By direct computation we have

$$\begin{aligned} & (\phi\psi)_{xxx} - \lambda(2\rho_x\phi\psi + 4\rho(\phi\psi)_x) \\ &= \phi_{xxx}\psi + 3\phi_{xx}\psi_x + 3\phi_x\psi_{xx} + \phi\psi_{xxx} - \lambda(4\rho\phi_x\psi + 4\rho\phi\psi_x + 2\rho_x\phi\psi) \\ &= (\phi_{xx} - \lambda\rho\phi)_x\psi + (\psi_{xx} - \lambda\rho\psi)_x\phi + 3(\phi_{xx} - \lambda\rho\phi)\psi_x + 3(\psi_{xx} - \lambda\rho\psi)\phi_x = 0, \end{aligned}$$

and also

$$W(\phi^2, \phi\psi, \psi^2) = 2W(\phi, \psi)^3,$$

where W denotes the Wronskian. The first computation shows that these indeed form solutions of the third order problem, while the second computation shows they span its solution space provided ϕ, ψ are linearly independent. \square

Lemma 5.2. Let $\lambda \in \mathbb{C}$, $\rho \in C^1([0, 1])$ and suppose $\rho(0) = \rho(1) = 0$. Moreover, let $\omega = \phi\psi$ where ϕ and ψ satisfy

$$\begin{aligned}\phi_{xx} &= \lambda\rho\phi, & \phi_x(0) - h\phi(0) &= 0, \\ \psi_{xx} &= \lambda\rho\psi, & \psi_x(1) + H\psi(1) &= 0.\end{aligned}$$

Then ω satisfies

$$\frac{1}{2}\omega_{xxx} - \lambda\mathcal{L}_\rho\omega = 0, \quad (5.2)$$

with the boundary conditions (3.3):

$$\begin{aligned}\frac{1}{2}\omega_{xx}(0) - h\omega_x(0) + h^2\omega(0) &= 0, \\ \frac{1}{2}\omega_{xx}(1) + H\omega_x(1) + H^2\omega(1) &= 0.\end{aligned}$$

Proof. That such a ω satisfies the third order differential equation is guaranteed by Lemma 5.1. To check the boundary conditions we calculate them explicitly. In particular, $\omega_{xx} = 2\phi_x\psi_x$ at $x = 0, 1$ because $\rho = 0$ at $x = 0, 1$. Then at $x = 0$ (3.3) becomes

$$(\phi_x\psi_x - h\phi_x\psi - h\phi\psi_x + h^2\phi\psi)|_{x=0} = (\psi_x - h\psi)(\phi_x - h\phi)|_{x=0} = 0,$$

and similarly at $x = 1$

$$(\phi_x\psi_x + H\phi_x\psi + H\phi\psi_x + H^2\phi\psi)|_{x=1} = (\phi_x + H\phi)(\psi_x + H\psi)|_{x=1} = 0.$$

□

Our attention will now turn to the determination of β (see (4.2)) and, subsequently, K from Proposition 4.1.

First we recall the definition of the *bilinear concomitant* ([17], p.125) associated with the third order differential equation $y''' = 0$.

Definition 5.3. Let $f, g \in C^3$ then the bilinear concomitant is

$$B(f, g)(x) = f_{xx}g - f_xg_x + fg_{xx}.$$

The second pertinent fact is the skew-symmetry of the operator \mathcal{L}_ρ relative to the natural inner product $(f, g) = \int_0^1 fgd x$. Indeed, whenever $\rho(0) = \rho(1) = 0$ as one easily checks using integration by parts:

$$(f, \mathcal{L}_\rho g) = -(\mathcal{L}_\rho f, g)$$

holds.

To proceed further we establish two important lemmas.

Lemma 5.4. Let $\lambda \in \mathbb{C}$, $\rho \in C^1([0, 1])$ and suppose $\rho(0) = \rho(1) = 0$. If ω satisfies $\frac{1}{2}\omega_{xxx} - \lambda\mathcal{L}_\rho\omega = 0$ with boundary conditions (3.3) then

$$\begin{aligned} \left(\frac{1}{2}\omega_x(0) - h\omega(0)\right)^2 &= \left(\frac{1}{2}\omega_x(0) - \frac{1}{2h}\omega_{xx}(0)\right)^2 \\ &= \left(\frac{1}{2}\omega_x(1) + H\omega(1)\right)^2 = \left(\frac{1}{2}\omega_x(1) + \frac{1}{2H}\omega_{xx}(1)\right)^2, \end{aligned} \quad (5.3)$$

where the first and the third equalities are displayed in order to deal with the cases of $h = \infty$, $H = \infty$ respectively.

Proof. Since ω satisfies $\frac{1}{2}\omega_{xxx} = \lambda\mathcal{L}_\rho\omega$ we easily obtain, using the skew-symmetry of \mathcal{L}_ρ , that $0 = (\omega, \omega_{xxx}) = \frac{1}{2}B(\omega, \omega)|_0^1$ which is equivalent to

$$\frac{1}{2}\omega_x^2(1) - \omega(1)\omega_{xx}(1) = \frac{1}{2}\omega_x^2(0) - \omega(0)\omega_{xx}(0). \quad (5.4)$$

Next we solve (3.3) for $\omega(j)$ or $\omega_{xx}(j)$, with $j = 0, 1$, depending on whether we want to consider finite or infinite h or H , subsequently obtaining

$$\omega(0) = \frac{1}{h}\omega_x(0) - \frac{1}{2h^2}\omega_{xx}(0), \quad \omega_{xx}(0) = 2h\omega_x(0) - 2h^2\omega(0),$$

and likewise at $x = 1$ with h replaced by $-H$. The result follows upon substituting these expressions back into (5.4) and simplifying. \square

Lemma 5.5. Let ω be given as in Lemma 5.2 and suppose $\lambda > 0$. Then

- (1) $\frac{1}{2}\omega_x(0) - h\omega(0) \neq 0$, for $0 \leq h < \infty$ or $\frac{1}{2}\omega_x(0) - \frac{1}{2h}\omega_{xx}(0) \neq 0$, for $0 < h \leq \infty$;
- (2) $\frac{1}{2}\omega_x(1) + H\omega(1) \neq 0$, for $0 \leq H < \infty$ or $\frac{1}{2}\omega_x(1) + \frac{1}{2H}\omega_{xx}(1) \neq 0$, for $0 < H \leq \infty$.

If $\lambda = 0$ inequalities (1) and (2) hold with the exception of Neumann-Neumann boundary conditions $H = h = 0$ in which case $\omega_x(0) = \omega_x(1) = 0$.

Proof. First, let us consider the second statement of item (1). Since by Lemma 5.2 $\omega = \phi\psi$ we can write

$$\begin{aligned} \frac{1}{2}\omega_x(0) - \frac{1}{2h}\omega_{xx}(0) &= \left(\frac{1}{2}\phi_x\psi + \frac{1}{2}\phi\psi_x - \frac{1}{h}\phi_x\psi_x\right)|_{x=0} \\ &= \left(\frac{\phi_x}{2}\left(\psi - \frac{1}{h}\psi_x\right) + \frac{\psi_x}{2}\left(\phi - \frac{1}{h}\phi_x\right)\right)|_{x=0} = \left(\frac{\phi_x}{2}\left(\psi - \frac{1}{h}\psi_x\right)\right)|_{x=0}. \end{aligned} \quad (5.5)$$

Thus it remains to show that both $\phi_x(0) \neq 0$ and $\psi(0) - \frac{1}{h}\psi_x(0) \neq 0$. Indeed, $\phi_x(0) \neq 0$ since $\phi(0) - \frac{1}{h}\phi_x(0) = 0$, and $h > 0$ means that if $\phi_x(0) = 0$ then also $\phi(0) = 0$ which would yield only trivial solutions. Now to show $\psi(0) - \frac{1}{h}\psi_x(0) \neq 0$ we multiply the corresponding second order problem by ψ and integrate from 0 to 1

$$\int_0^1 \psi\psi_{xx} dx = \lambda \int_0^1 \rho\psi^2 dx \iff \psi\psi_x|_0^1 - \int_0^1 \psi_x^2 dx = \lambda \int_0^1 \rho\psi^2 dx.$$

Solving for λ and substituting in the boundary condition on the right gives

$$\lambda = -\frac{H\psi(1)^2 + \psi(0)\psi_x(0) + \int_0^1 \psi_x^2 dx}{\int_0^1 \rho\psi^2 dx},$$

from which we see that if $\psi(0) - \frac{1}{h}\psi_x(0) = 0$ then

$$\lambda = -\frac{H\psi(1)^2 + \frac{1}{h}\psi_x(0)^2 + \int_0^1 \psi_x^2 dx}{\int_0^1 \rho\psi^2 dx} \leq 0,$$

contradicting our our assumption $0 < \lambda$.

For the first statement of item (1), we begin by computing

$$\frac{1}{2}\omega_x(0) - h\omega(0) = \frac{\phi(0)}{2}(\psi_x(0) - h\psi(0)), \quad (5.6)$$

and then proceed in a similar way to the argument in support of the second statement of item (1).

For item (2) we find that

$$\begin{aligned} \frac{1}{2}\omega_x(1) + \frac{1}{2H}\omega_{xx}(1) &= \frac{\psi_x(1)}{2}(\phi(1) + \frac{1}{H}\phi_x(1)), \\ \frac{1}{2}\omega_x(1) + H\omega(1) &= \frac{\psi(1)}{2}(\phi_x(1) + H\phi(1)), \end{aligned} \quad (5.7)$$

from which the same conclusions as in the last two cases can be made by similar arguments. Finally the case $\lambda = 0$ can be checked by direct computation using $\phi(x) = hx + 1, \psi(x) = H(1 - x) + 1$. \square

Lemma 5.6. *Let $\omega = \omega(x; \lambda)$ be given as in Lemma 5.2. Then*

1. ω is entire in λ of order $\frac{1}{2}$;
2. the following equalities hold (whenever they make sense)

$$\begin{aligned} \left(\frac{1}{2}\omega_x(0) - \frac{1}{2h}\omega_{xx}(0)\right) &= -\left(\frac{1}{2}\omega_x(0) - h\omega(0)\right) \\ &= -\left(\frac{1}{2}\omega_x(1) + \frac{1}{2H}\omega_{xx}(1)\right) = \left(\frac{1}{2}\omega_x(1) + H\omega(1)\right). \end{aligned} \quad (5.8)$$

Proof. (1) By Lemma 5.2 it suffices to show that the solution to $\phi_{xx} - \lambda\rho\phi = 0$ with $\phi_x(0) - h\phi(0) = 0$ or $\phi_x(1) + H\phi(1) = 0$ is analytic in λ . This part is well known but we will need some intermediate formulas so, for completeness, we proceed with the proof, referring to the monograph [20] for details. We begin by writing ϕ as a power series in λ :

$$\phi(x; \lambda) = c_0(x) + \sum_{n=1}^{\infty} \lambda^n c_n(x),$$

which when substituted into the second order differential equation yields the system

$$c_{0,xx} = 0, \quad c_{n,xx} - c_{n-1}\rho = 0 \quad (n \geq 1).$$

Since the argument proceeds in the same way regardless of whether we use the left or right initial condition we concentrate now on the left initial condition.

We immediately see that c_0 will be linear in x with two unknowns which can be chosen to satisfy the initial condition on the left. Then we choose c_n so that $c_n(0) = c_{n,x}(0) = 0$, holds for $n \geq 1$, ensuring that ϕ satisfies the left initial condition. Once this is accomplished we must show that the power series converges.

To determine c_n ($n \geq 1$) we simply integrate the c_n relation from 0 to x . Then by recursion we find that

$$c_n(x) = \int_{0 \leq \xi_1 \leq \xi_2 \leq \dots \leq \xi_{n+1} = x} \left[\prod_{j=1}^n (\xi_{j+1} - \xi_j) \rho(\xi_j) \right] c_0(\xi_1) d\xi_1 d\xi_2 \dots d\xi_n \quad (5.9)$$

To show that the power series converges it suffices to observe that

$$\begin{aligned} |c_n(x)| &\leq C x^n \int_{0 \leq \xi_1 \leq \xi_2 \leq \dots \leq \xi_{n+1} = x} \left[\prod_{j=1}^n \rho(\xi_j) \right] d\xi_1 d\xi_2 \dots d\xi_n \\ &\leq C \frac{x^n}{n!} \left(\int_{[0,x]} \rho(\xi) d\xi \right)^n \stackrel{\text{def}}{=} C \frac{x^n}{n!} (M(x))^n, \end{aligned} \quad (5.10)$$

where $|c_0(x)| < C$ and $M(x)$ is the total mass on the interval $[0, x]$. Hence, the power series $\phi(x; \lambda)$ converges uniformly on compact sets in $\lambda \in \mathbb{C}$. One can improve on (5.10) by estimating more accurately the term $(x - \xi_n) \prod_{j=1}^{n-1} (\xi_{j+1} - \xi_j)$ in the region $\{0 \leq \xi_1 \leq \xi_2 \leq \dots \leq \xi_n \leq x\}$. For example, one can check by induction that

$$(x - \xi_n) \prod_{j=1}^{n-1} (\xi_{j+1} - \xi_j) \leq \left(\frac{x}{n} \right)^n$$

and, upon replacing x^n in (5.10) with this more accurate estimate, we obtain

$$|c_n(x)| \leq C \frac{(xM(x))^n}{n!n^n}.$$

Moreover, since $2^n n^n n! > (2n)!$, one easily obtains

$$|\phi(x; \lambda)| \leq C \cosh \left(\sqrt{2xM(x)|\lambda|} \right),$$

implying that $\phi(x; \lambda)$ is an entire function of order not exceeding $\frac{1}{2}$. That the order is exactly $\frac{1}{2}$ requires more subtle analysis, for example of the growth of zeros of ϕ , which can be found in M.G. Krein's work [19] and citations therein.

Since the computation for $\psi(x; \lambda)$ is similar we record only the main steps:

$$\begin{aligned} \psi(x; \lambda) &= \hat{c}_0(x) + \sum_{n=1} \lambda^n \hat{c}_n(x) \\ \hat{c}_{0,xx} &= 0, \quad \hat{c}_{n,xx} = \rho \hat{c}_{n-1}, \quad \hat{c}_{0,x}(1) + H\hat{c}_0(1) = 0, \quad \hat{c}_{n,x}(1) = \hat{c}_n(1) = 0 \\ \hat{c}_n &= \int_{0 \leq \xi_{n+1} = x \leq \xi_n \leq \dots \leq \xi_1 \leq 1} \left[\prod_{j=1}^n (\xi_j - \xi_{j+1}) \rho(\xi_j) \right] \hat{c}_0(\xi_1) d\xi_1 d\xi_2 \dots d\xi_n. \end{aligned} \quad (5.11)$$

Finally the growth estimate reads:

$$|\hat{c}_n|(x) \leq \hat{C} \frac{((1-x)\hat{M}(x))^n}{n!n^n},$$

where $\hat{C}, \hat{M}(x)$ denote the bound on \hat{c}_0 , the mass of $[x, 1]$ respectively. Thus ψ is also entire of order not exceeding $\frac{1}{2}$, hence so is the product of ϕ and ψ .

- (2) Except for the special case $h = H = 0$ (i.e. Neumann-Neumann boundary conditions) (see Lemma 5.5) we proceed by showing first that the equalities hold when $\lambda = 0$ and extend them by continuity to a neighbourhood of $\lambda = 0$. Then the result for $\lambda > 0$ follows by using Lemmas 5.4, 5.5 and analyticity of b as a function of λ proved in item (1).

Indeed, for $\lambda = 0$ and we can take $c_0 = hx + 1, \hat{c}_0 = H(1-x) + 1, \omega = c_0\hat{c}_0$, and use equations (5.6), (5.7) to verify the equalities. For boundary conditions other than Neumann-Neumann all quantities in brackets involving b are nonzero so we can form ratios. By Lemma 5.4 all these ratios squared are 1 for any λ . Since these ratios are -1 for $\lambda = 0$ they have to stay -1 in a neighbourhood of $\lambda = 0$ by continuity. Thus equations (5.8) hold in the whole domain of analyticity (C) by the Identity Theorem.

For the special case $h = H = 0$ it suffices to check that the equalities hold in a neighbourhood of $\lambda = 0$ by examining ω up to the first order in λ . Since in this case $c_0 = \hat{c}_0 = 1$ we obtain

$$\begin{aligned} \omega &= 1 + (c_1 + \hat{c}_1)\lambda + \mathcal{O}(\lambda^2) = 1 + \left(\int_0^x (x-\xi)\rho(\xi)d\xi + \int_x^1 (\xi-x)\rho(\xi)d\xi \right)\lambda + \mathcal{O}(\lambda^2), \\ \omega_x &= \left(\int_0^x \rho(\xi)d\xi - \int_x^1 \rho(\xi)d\xi \right)\lambda + \mathcal{O}(\lambda^2), \end{aligned}$$

which, when evaluated, give $\omega_x(0) = -(\int_0^1 \rho(\xi)d\xi)\lambda + \mathcal{O}(\lambda^2), \omega_x(1) = (\int_0^1 \rho(\xi)d\xi)\lambda + \mathcal{O}(\lambda^2)$, hence proving (5.8) for $h = 0 = H$. □

Recall that the eigenvalues of the original boundary value problem (1.2) are roots of $v_x(1; z) + Hv(1; z) \stackrel{\text{def}}{=} D(z)$ which we can write either in terms of ϕ or ψ , after adjusting the overall multiplier, for example by agreeing that $v(x; z) = hx + 1 = c_0$ on the massless portion to the right of 0. Taking into account Lemma 5.1 we easily obtain:

$$D(-\lambda) = \phi_x(1; \lambda) + H\phi(1; \lambda) = -(\psi_x(0; \lambda) - h\psi(0; \lambda)).$$

One useful way of computing D is by using the expansion in terms of iterates c_n (see equation (5.9)).

Lemma 5.7. *Let $\rho \in C^1([0, 1])$ and let $W(f, g) = fg' - f'g$ denote the Wronskian of two functions f, g . Suppose that $h + H > 0$ and let the eigenvalues of the string problem (1.2) be denoted by $0 < z_1 < z_2 < \dots$, then D admits the additive representation*

$$D(-\lambda) = W(\hat{c}_0, c_0) + \sum_{n=1}^{\infty} \lambda^n \int_{0 \leq \xi_1 \leq \xi_2 \leq \dots \leq \xi_n \leq 1} \hat{c}_0(\xi_n)\rho(\xi_n) \left(\prod_{j=1}^{n-1} (\xi_{j+1} - \xi_j)\rho(\xi_j) \right) c_0(\xi_1) d\xi_1 d\xi_2 \dots d\xi_n, \quad (5.12)$$

and the product representation

$$D(-\lambda) = W(\hat{c}_0, c_0) \prod_{j=1}^{\infty} \left(1 + \frac{\lambda}{z_j}\right). \quad (5.13)$$

Proof. The additive formula follows from the definitions (5.9), (5.11) of c_n, \hat{c}_n , respectively. The multiplicative formula is well known and it follows immediately from two facts: ϕ is an entire function of order $\frac{1}{2}$ and $\frac{z_n}{n^2} = \mathcal{O}(1)$, $n \rightarrow \infty$, and it therefore admits Hadamard's product formula. \square

Now we are ready to give a complete construction of b_{-1} and b_0 appearing in equation (5.1). The notation in the Theorem below highlights the dependence on the parameter ϵ ; for more details regarding this notation we refer to the proof of Lemma 5.6.

Theorem 5.8. *Let $\rho \in C^1([0, 1])$, β, K be defined as in Proposition 4.1, using equations (4.2), (4.6) respectively, and let ω be given as in Lemma 5.2. Let*

$$\begin{aligned} b_{-1}(x; \epsilon) &= \omega(x; \epsilon), \\ b_0(x; \epsilon) &= \frac{[\omega(x; 0) - \omega(x; \epsilon)]}{\epsilon}, \end{aligned}$$

then b_{-1} and b_0 so defined satisfy (5.1b).

Moreover, for this choice of b_{-1}, b_0 ,

1. $\beta(z; \epsilon) = \frac{1}{2} \left[\frac{(D(-\epsilon) - D(0))}{\epsilon} - \frac{D(-\epsilon)}{z + \epsilon} \right]$,
2. $K = -1$,
3. $\beta(z; \epsilon)$ is deformation invariant ($\dot{\beta} = 0$).

Proof. The formula for b_{-1} follows from Lemma 5.2 for $\lambda = \epsilon$ to conform with equations (5.1). Likewise, the formula for b_0 is a solution of the second equation in (5.1). The additional term involving $c_0 \hat{c}_0$ is added to ensure the existence of the limit $\epsilon \rightarrow 0$. To prove the formula for β one needs to consider all individual cases.

1. $0 \leq h < \infty$: we use the formula (5.6) which we recall for the sake of clarity, explicitly displaying the ϵ dependence $\frac{1}{2} b_{-1,x}(0; \epsilon) - h b_{-1}(0; \epsilon) = \frac{1}{2} \phi(0; \epsilon) (\psi_x(0; \epsilon) - h \psi(0; \epsilon)) = -\frac{1}{2} D(-\epsilon)$ where we used that $\phi(0; \epsilon) = 1$ by construction. For b_0 , using linearity, we obtain $\frac{1}{2} b_{0,x}(0; \epsilon) - h b_0(0; \epsilon) = \frac{1}{2\epsilon} (D(-\epsilon) - D(0))$
2. $0 < h \leq \infty$: to compute β in this case we use the second formula in equation 4.2 as well as (5.5) to obtain $-\frac{1}{2} b_{-1,x}(0; \epsilon) + \frac{1}{2h} b_{-1,xx}(0; \epsilon) = -\frac{1}{2} D(-\epsilon)$, recalling that for this range of h , $c_0 = x + \frac{1}{h}$; similar argument to the one used in the first item works for b_0 .

The computation of β follows now from the definition (4.2). The statement that β is constant under the deformation, in other words β is a constant of "motion", is an automatic consequence of isospectrality which implies that $\dot{D}(\lambda) = 0$; in principle there could be an overall time-dependent factor but by construction the coefficient of λ^0 in $D(\lambda)$ is $(c_{0,x} + H c_0)(1)$ which is constant in t . Finally, $K = -1$ (see (4.6)) follows from (5.8). \square

We conclude this section with an appropriate modification of Theorem 5.8 for the rational flows (2.8). The proof is a straightforward generalization of the previous proof, taking into account that the fields b_0 and b_k in equation (2.9b) can be rescaled.

Theorem 5.9. *Let $\rho \in C^1([0, 1])$, β, K be defined as in Proposition 4.1 using equations (4.2), (4.6) respectively, and let ω be given as in Lemma 5.2. Given M distinct positive numbers $\epsilon_k, 1 \leq k \leq M$, and $M+1$ functions of $\epsilon = (\epsilon_1, \dots, \epsilon_M)$, denoted $\mu_0, \mu_1, \dots, \mu_M$, let*

$$b_{-1}^{(k)}(x) = \mu_k \omega(x; \epsilon_k), \quad 1 \leq k \leq M$$

$$b_0(x) = \sum_{k=1}^M \frac{[\mu_0 \omega(x; 0) - \mu_k \omega(x; \epsilon_k)]}{\epsilon_k}.$$

Then $b_{-1}^{(k)}$ and b_0 so defined satisfy

$$b_{0,xxx} + \sum_{k=1}^M \frac{b_{-1,xxx}^{(k)}}{\epsilon_k} = 0, \quad \frac{1}{2} b_{-1,xxx}^{(k)} - \epsilon_k \mathcal{L}_\rho b_{-1}^{(k)} = 0, \quad k = 1 \dots M,$$

(see (2.9b)). Moreover, for this choice of $b_{-1}^{(k)}, b_0$, setting $b = b_0 + \sum_{k=1}^M \frac{b_{-1}^{(k)}}{z + \epsilon_k}$, one obtains

1.

$$\beta(z; \epsilon) = \frac{1}{2} \sum_{k=1}^M \left[\frac{(\mu_k D(-\epsilon_k) - \mu_0 D(0))}{\epsilon_k} - \mu_k \frac{D(-\epsilon_k)}{z + \epsilon_k} \right], \quad (5.14)$$

2. $K = -1$,

3. $\beta(z; \epsilon)$ is deformation invariant.

5.1 Discrete mass density

In this short section we will discuss necessary modifications of the formalism for the non-smooth case of

$$\rho = \sum_{j=1}^N m_j \delta_{x_j}, \quad 0 < x_1 < \dots < x_N < 1, \quad (5.15)$$

with weights (point masses) $m_j > 0$. We will leave the details of a general non-smooth case to future work; our goal is to argue that the formalism is valid for the case of discrete positive measures with only minor modifications. We start with the discussion of the boundary value problem (1.2). The only modification needed to deal with discrete measures is to take derivatives as distributional derivatives. Thus the boundary value problem simply reads:

$$-D_x^2 v = z\rho(x)v, \quad 0 < x < 1, \quad v_x(0) - hv(0) = 0, \quad v_x(1) + Hv(1) = 0, \quad (5.16)$$

implying that in the case of the discrete ρ , v will be continuous and piecewise differentiable, with jumps on the support of ρ . For such functions the distributional derivatives can be conveniently computed from

$$D_x f = f_x + \sum_{j=1}^N [f](x_j) \delta_{x_j}, \quad (5.17)$$

and its higher order extensions; in the formula (5.17) f_x denotes the classical derivative, $[f](x_j)$ means the jump of f at x_j . Thus, away from the support of ρ , the equation reads: $-v_{xx} = 0$, while on the support of ρ we obtain:

$$-[v_x](x_j) = z m_j v(x_j).$$

We will also need to extend the action of \mathcal{L}_ρ to continuous, piecewise differentiable functions f :

$$\mathcal{L}_\rho f \stackrel{\text{def}}{=} D_x(f\rho) + \langle f_x \rangle \rho = D_x(f\rho) + \sum_{j=1}^N \langle f_x \rangle(x_j) m_j \delta_{x_j} \quad (5.18)$$

where $\langle f \rangle$ is an everywhere defined average function of pointwise right-hand and left-hand limits.

We will now examine appropriate modifications of Lemmas 5.2, 5.1 which are critical for the whole formalism.

Lemma 5.10. *Let $\lambda \in \mathbb{C}$, ρ be a discrete measure given by (5.15) and suppose ϕ and ψ are solutions of the distributional equation $D_x^2 f - \lambda \rho f = 0$. Then $\phi\psi$ is a solution of the distributional equation $\frac{1}{2}D_x^3 g - \lambda \mathcal{L}_\rho g = 0$.*

Proof. We observe that both ϕ and ψ are continuous and piecewise differentiable; in fact piecewise linear. Away from the support of ρ the claim holds by the original Lemma 5.1. On the support, that is at the points x_j , one needs to compute the contributions of distributional derivatives. One can localize the problem by choosing test functions to be supported only in the neighbourhood of individual points of the support. Hence, we can perform computations as if ρ were $\rho = m_1 \delta_{x_1}$, thereby dropping the sum from the computation, only to reinstate it in the final answer. We compute now $D_x^3(\phi\psi)$:

$$\begin{aligned} D_x^3(\phi\psi) &= D_x^2((\phi\psi)_x) = D_x((\phi\psi)_{xx} + [(\phi\psi)_x](x_1)\delta_{x_1}) = \\ &(\phi\psi)_{xxx} + [(\phi\psi)_x](x_1)\delta'_{x_1} + [(\phi\psi)_{xx}](x_1)\delta_{x_1} = [(\phi\psi)_x](x_1)\delta'_{x_1} + [(\phi\psi)_{xx}](x_1)\delta_{x_1}, \end{aligned}$$

where we used that ϕ, ψ are piecewise linear (in x) continuous functions. We recall that $[\phi_x](x_1) = \lambda m_1 \phi(x_1)$, $[\psi_x](x_1) = \lambda m_1 \psi(x_1)$, which implies

$$[(\phi\psi)_x](x_1)\delta'_{x_1} = 2\lambda\phi(x_1)\psi(x_1)m_1\delta'_{x_1} = 2\lambda D_x((\phi\psi)\rho).$$

Likewise, for the second singular term we obtain:

$$\begin{aligned} [(\phi\psi)_{xx}](x_1)\delta_{x_1} &= 2[\phi_x\psi_x](x_1)\delta_{x_1} = 2([\phi_x](x_1)\langle\psi_x\rangle(x_1) + [\psi_x](x_1)\langle\phi_x\rangle(x_1))\delta_{x_1} = \\ &2(\lambda m_1 \phi(x_1)\langle\psi_x\rangle(x_1) + \lambda m_1 \psi(x_1)\langle\phi_x\rangle(x_1))\delta_{x_1} = 2\lambda\langle(\phi\psi)_x\rangle\rho, \end{aligned}$$

which, in conjuncture with the first calculation and in view of equation (5.18), implies

$$\frac{1}{2}D_x^3(\phi\psi) = \lambda \mathcal{L}_\rho(\phi\psi),$$

which completes the proof. □

Lemma 5.11. *Let $\lambda \in \mathbb{C}$ and ρ be the discrete measure defined by (5.15). Moreover, let $\omega = \phi\psi$ where ϕ and ψ satisfy*

$$\begin{aligned} D_x^2 \phi &= \lambda \rho \phi, & \phi_x(0) - h\phi(0) &= 0, \\ D_x^2 \psi &= \lambda \rho \psi, & \psi_x(1) + H\psi(1) &= 0. \end{aligned}$$

Then ω satisfies

$$\frac{1}{2} D_x^3 \omega - \lambda \mathcal{L}_\rho \omega = 0, \quad (5.19)$$

with the boundary conditions (3.3):

$$\begin{aligned} \frac{1}{2} \omega_{xx}(0) - h\omega_x(0) + h^2 \omega(0) &= 0, \\ \frac{1}{2} \omega_{xx}(1) + H\omega_x(1) + H^2 \omega(1) &= 0. \end{aligned}$$

Proof. The only thing to check are the boundary conditions. However, the support of ρ is away from the end points so the computation takes place in the smooth sector and, as a consequence, the computation done previously applies. \square

Finally, we present appropriate analogs of Lemma 5.7 and Theorem 5.8. We recall that the function defining the spectrum of the boundary value problem (1.2) is given, in our notation, by $D(-\lambda) = \phi_x(1; \lambda) + H\phi(1; \lambda)$, $0 \leq H$, where the eigenvalue z in the original problem is taken to be $-\lambda$. We also recall (see (5.9), (5.11)) that $c_0(x) = \phi(x; 0)$, $\hat{c}_0(x) = \psi(x; 0)$.

Lemma 5.12. *Let ρ be the discrete measure defined by (5.15) and let $W(f, g)$ denote the Wronskian of two functions f, g . Suppose that $h + H > 0$ and let the eigenvalues of the string problem (1.2) be denoted by $0 < z_1 < z_2 < \dots < z_N$, then D admits the additive representation*

$$D(-\lambda) = W(\hat{c}_0, c_0) + \sum_{n=1}^N \lambda^n \sum_{1 \leq i_1 < i_2 < \dots < i_n \leq N} m_{i_1} m_{i_2} \dots m_{i_n} (x_{i_n} - x_{i_{n-1}})(x_{i_{n-1}} - x_{i_{n-2}}) \dots (x_{i_2} - x_{i_1}) c_0(x_{i_1}) \hat{c}_0(x_{i_n}), \quad (5.20)$$

and the product representation

$$D(-\lambda) = W(\hat{c}_0, c_0) \prod_{j=1}^N \left(1 + \frac{\lambda}{z_j} \right). \quad (5.21)$$

Proof. The additive formula follows from the definitions (5.9), (5.11) of c_n, \hat{c}_n , respectively. The multiplicative formula is well known and it follows immediately from Hadmard's product formula and the fact that the spectrum is finite because m is a finite discrete measure. \square

Theorem 5.13. *Let ρ be the discrete measure given by (5.15), β, K be defined as in Proposition 4.1, using equations (4.2), (4.6) respectively, and let ω be given as in Lemma 5.11. Then*

1. ϕ, ψ defined in Lemma 5.11 are polynomials of degree N ;
2. $D(-\lambda) = \phi_x(1; \lambda) + H\phi(1; \lambda) = -(\psi_x(0; \lambda) - h\psi(0; \lambda))$;

3. b_{-1}, b_0 defined by

$$\begin{aligned} b_{-1}(x; \epsilon) &= \omega(x; \epsilon), \\ b_0(x; \epsilon) &= \frac{[\omega(x; 0) - \omega(x; \epsilon)]}{\epsilon}, \end{aligned} \quad (5.22)$$

satisfy equation (5.1b) with \mathcal{L}_ρ defined by (5.18).

4.

$$\beta(z; \epsilon) = \frac{1}{2} \left[\frac{D(-\epsilon) - D(0)}{\epsilon} - \frac{D(-\epsilon)}{z + \epsilon} \right],$$

5. $K = -1$, and $\beta(z; \epsilon)$ is deformation invariant ($\dot{\beta} = 0$).

6. Moreover, if $h + H > 0$ and $0 < z_1 < z_2 < \dots < z_N$ denotes the spectrum of the boundary value problem (1.2) then

$$D(-\epsilon) = W(\hat{c}_0, c_0) \prod_{j=1}^N \left(1 + \frac{\epsilon}{z_j} \right). \quad (5.23)$$

Proof. The assertions of the theorem follow from the fact that c_n and \hat{c}_n defined by equations (5.9) and (5.11) are valid if ρ is any positive measure, in particular if ρ is finite and discrete. The series (in λ) which define ϕ and ψ (see the proof of Lemma 5.6) terminate at $n = N$ as can be easily verified. For example: the term $\prod_{j=1}^n (\xi_{j+1} - \xi_j) \rho(\xi_j)$ will vanish identically on the domain $0 \leq \xi_1 \leq \xi_2 \leq \dots \leq \xi_{n+1} = x$ if $n > N$. In particular $D(z) = \phi_x(1; -z) + H\phi(1; -z)$ is a polynomial of degree N and $D(z) = -(\psi_x(0; -z) - h\psi(0; -z))$ by either directly inspecting iterates c_n and \hat{c}_n or by observing that $\phi_x(1; \lambda) + H\phi(1; \lambda)$ and $\psi_x(0; \lambda) - h\psi(0; \lambda)$ can only differ by a constant multiplier, independent of λ , which can easily be verified to be -1 for $\lambda = 0$. The statements about β and K then follow from equations (5.5), (5.6), (5.7). Finally, the invariance of the spectrum implies that β is constant in t . \square

We recall the deformation equation (2.13)

$$zD_t \rho = \frac{1}{2} D_x^3 b + z \mathcal{L}_\rho b.$$

Evaluating this equation on the support of ρ with the help of the previous theorem and the definition of \mathcal{L}_ρ given by (5.18) we obtain the following explicit form of the evolution equation, which we state for simplicity in the simplest case of a single pole rational flow.

Corollary 5.14. *Let the ZS deformation be given by $b = b_0 + \frac{b_{-1}}{z+\epsilon}$, $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, $0 < x_1 < \dots < x_N < 1$, and let b_0 and $b_{-1}^{(k)}$ be given by equations (5.22). Then the positions x_j and masses m_j evolve according to:*

$$\boxed{\dot{x}_j = -b_0(x_j), \quad \dot{m}_j = m_j \langle b_{0,x} \rangle(x_j), \quad 1 \leq j \leq N.}$$

The sums

$$\boxed{I_1 = \sum_{i=1}^N m_i c_0(x_i) \hat{c}_0(x_i),}$$

and

$$I_n = \sum_{1 \leq i_1 < i_2 < \dots < i_n \leq N} m_{i_1} m_{i_2} \dots m_{i_n} (x_{i_n} - x_{i_{n-1}})(x_{i_{n-1}} - x_{i_{n-2}}) \dots (x_{i_2} - x_{i_1}) c_0(x_{i_1}) \hat{c}_0(x_{i_n}), \quad 2 \leq n \leq N$$

are invariants of the ZS deformation (constants of motion).

It is worth mentioning that the deformation parameter t has no direct physical meaning, and as such it can be chosen arbitrarily. Since the constraints on the fields b_{-1}, b_0 are homogenous, one can rescale both fields by a common, x -independent factor. The rescaling by the Wronskian $W(c_0, \hat{c}_0)$ is particularly natural as it allows one to formulate equations in terms of Green's functions.

Corollary 5.15. *Let the ZS deformation be given by $b = b_0 + \frac{b_{-1}}{z+\epsilon}$ with $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, where $0 < x_1 < \dots < x_N < 1$, and let $h + H > 0$ in the boundary value problem (1.2). Let $G(x, y; \lambda)$ be Green's function of $D_x^2 - \lambda \rho$ to boundary conditions $G_x(0, y) - hG(0, y) = 0$, $G_x(1, y) + HG(1, y) = 0$ and define new b_0 and b_{-1} by:*

$$b_{-1}(x; \epsilon) \stackrel{\text{def}}{=} G(x, x; \epsilon), \quad b_0(x; \epsilon) \stackrel{\text{def}}{=} \frac{1}{\epsilon} (G(x, x; 0) - G(x, x; \epsilon)).$$

Then the positions x_j and masses m_j evolve with respect to the rescaled time $\tilde{t} = W(c_0, \hat{c}_0)t$ according to:

$$\dot{x}_j = -b_0(x_j), \quad \dot{m}_j = m_j \langle b_{0,x} \rangle(x_j), \quad 1 \leq j \leq N,$$

and the sums

$$I_1 = \sum_{i=1}^N m_i G(x_i, x_i; 0),$$

and

$$I_n = \sum_{1 \leq i_1 < i_2 < \dots < i_n \leq N} m_{i_1} m_{i_2} \dots m_{i_n} (x_{i_n} - x_{i_{n-1}})(x_{i_{n-1}} - x_{i_{n-2}}) \dots (x_{i_2} - x_{i_1}) G(x_{i_1}, x_{i_n}; 0), \quad 2 \leq n \leq N$$

are invariants of the ZS deformation (constants of motion).

5.2 Linearization of the spectral data

Using the expression for β from Theorems 5.8, 5.13 along with Proposition 4.1, in particular equation (4.6), we find that the following statements can be made for general boundary conditions. For brevity we will use the abbreviations $\mathcal{D}, \mathcal{N}, \mathcal{M}$ for Dirichlet, Neumann, or mixed boundary conditions respectively as well we will suppress, till further notice, the dependence on x, z and ϵ .

Theorem 5.16. *Let the ZS deformation be given by $b = b_0 + \frac{b_{-1}}{\epsilon+z}$, $\epsilon > 0$ with the mass density $\rho \in C^1([0, 1])$ which vanishes in a neighbourhood of each endpoint, or $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, $0 < x_1 < \dots < x_N < 1$. Then the string boundary value problem (1.2) is isospectral with the mass density ρ evolving according to*

$$\begin{aligned} \rho_t &= \mathcal{L}_\rho b_0, \\ \epsilon b_{0,xxx} + b_{-1,xxx} &= 0, \\ \frac{1}{2} b_{-1,xxx} - \epsilon \mathcal{L}_\rho b_{-1,x} &= 0, \end{aligned}$$

with b_0, b_1, β constructed in Theorems 5.8, 5.13 and \mathcal{L}_ρ defined by equations (2.6), (5.18). In addition, the evolution of the spectral data A_1, B_1 (equation (4.3)) for general boundary conditions linearizes as follows:

1. $\mathcal{D} - \mathcal{D}$: the spectrum is given by $B_1 = 0$, the evolution by

$$\dot{A}_1 = 2\beta A_1 - \frac{1}{2} b_{xx}(1) B_1, \quad \dot{B}_1 = 0.$$

2. $\mathcal{D} - \mathcal{N}$: the spectrum is given by $A_1 = 0$, the evolution by

$$\dot{A}_1 = 0, \quad \dot{B}_1 = 2\beta B_1 + b(1) A_1.$$

3. $\mathcal{N} - \mathcal{N}$: the spectrum is given by $A_1 = 0$, the evolution by

$$\dot{A}_1 = 2\beta A_1, \quad \dot{B}_1 = b(1) A_1.$$

4. $\mathcal{N} - \mathcal{D}$: the spectrum is given by $B_1 = 0$, the evolution by

$$\dot{A}_1 = 2\beta A_1 - \frac{1}{2} b_{xx}(1) B_1, \quad \dot{B}_1 = 0.$$

5. $\mathcal{D} - \mathcal{M}$: the spectrum is given by $A_1 + HB_1 = 0$, the evolution by

$$\dot{A}_1 + H\dot{B}_1 = 0, \quad \dot{A}_1 - H\dot{B}_1 = 2\beta[A_1 - HB_1] + b_x(1)[A_1 + HB_1].$$

6. $\mathcal{M} - \mathcal{D}$: the spectrum is given by $B_1 = 0$, the evolution by

$$\dot{A}_1 = 2\beta A_1 - \frac{1}{2} b_{xx}(1) B_1, \quad \dot{B}_1 = 0.$$

7. $\mathcal{N} - \mathcal{M}$: the spectrum is given by $A_1 + HB_1 = 0$, the evolution by

$$\dot{A}_1 + H\dot{B}_1 = 0, \quad \dot{A}_1 - H\dot{B}_1 = 2\beta[A_1 - HB_1] + b_x(1)[A_1 + HB_1].$$

8. $\mathcal{M} - \mathcal{N}$: the spectrum is given by $A_1 = 0$, the evolution by

$$\dot{A}_1 = 0, \quad \dot{B}_1 = 2\beta B_1 + b(1) A_1.$$

9. $\mathcal{M} - \mathcal{M}$: the spectrum is given by $A_1 + HB_1 = 0$, the evolution by

$$\dot{A}_1 + H\dot{B}_1 = 0, \quad \dot{A}_1 - H\dot{B}_1 = 2\beta[A_1 - HB_1] + b_x(1)[A_1 + HB_1].$$

Remark 5.17. Observe that even though the same type of boundary conditions on the right end of the string formally yields the same flow equations of the spectral data, the coefficient β is sensitive to the boundary conditions on the left end of the string, and so is the field b .

Remark 5.18. Theorem 5.16 can be readily generalized to the case of multi-pole rational flows (2.8) along the lines of Theorem 5.9.

5.3 The limit $\epsilon \rightarrow 0$

In previous sections we constructed the fields $b_{-1}(x; \epsilon)$, $b_0(x; \epsilon)$ in such a way that both are entire functions of ϵ . We recall equations that ρ and b satisfy:

$$\rho_t = \mathcal{L}_\rho b_0, \quad b_{0,xxx} + \frac{b_{-1,xxx}}{\epsilon} = 0, \quad \frac{1}{2} b_{-1,xxx} - \epsilon \mathcal{L}_\rho b_{-1} = 0,$$

which we write as

$$\rho_t = \mathcal{L}_\rho b_0, \quad \frac{1}{2} b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad \frac{1}{2} b_{-1,xxx} - \epsilon \mathcal{L}_\rho b_{-1} = 0.$$

We now take the limit $\epsilon \rightarrow 0$. Without giving a detailed account of intermediate elementary steps we summarize the results in the following theorem.

Theorem 5.19. *Let the ZS deformation be given by $b = b_0 + \frac{b_{-1}}{\epsilon+z}$, $\epsilon > 0$ with the mass density $\rho \in C^1([0, 1])$ which vanishes in a neighbourhood of each endpoint, or $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, $0 < x_1 < \dots < x_N < 1$. Let $b_{-1}(x; \epsilon)$ and $b_0(x; \epsilon)$ be given as in theorems 5.8 and 5.13. Then*

1.

$$\begin{aligned} \lim_{\epsilon \rightarrow 0} b_{-1}(x; \epsilon) &= c_0(x) \hat{c}_0(x), \\ \lim_{\epsilon \rightarrow 0} b_0(x; \epsilon) &= -(c_0(x) \hat{c}_1(x) + c_1(x) \hat{c}_0(x)), \\ D(0) &= W(\hat{c}_0, c_0), \quad D'(0) = \int_0^1 \hat{c}_0(\xi) \rho(\xi) c_0(\xi) d\xi, \end{aligned}$$

where $W(f, g)$ denotes the Wronskian of f, g .

2. If in addition $h + H > 0$ and the fields b_0 and b_{-1} are redefined as:

$$\begin{aligned} b_{-1}(x) &\stackrel{\text{def}}{=} \frac{c_0(x) \hat{c}_0(x)}{W(c_0, \hat{c}_0)}, \\ b_0(x) &\stackrel{\text{def}}{=} \frac{-(c_0(x) \hat{c}_1(x) + c_1(x) \hat{c}_0(x))}{W(c_0, \hat{c}_0)}, \end{aligned} \tag{5.24}$$

then

$$\boxed{b_{-1}(x) = G(x, x), \quad b_0(x) = - \int_0^1 |x - \xi| G(x, \xi) \rho(\xi) d\xi}$$

where $G(x, y)$ is Green's function of D_x^2 to boundary conditions $G_x(0, y) - hG(0, y) = 0$, $G_x(1, y) + HG(1, y) = 0$ and $\beta(z)$ (see Theorem 5.16) satisfies

$$\boxed{\beta(z) = - \frac{1}{2} \left(\int_0^1 G(\xi, \xi) \rho(\xi) d\xi - \frac{1}{z} \right)}$$

3. With b_0, b_{-1} so defined ρ undergoes an isospectral deformation governed by

$$\rho_{\tilde{t}} = \mathcal{L}_\rho b_0, \quad \frac{1}{2} b_{0,xxx} + \mathcal{L}_\rho b_{-1} = 0, \quad b_{-1,xxx} = 0, \quad \text{where } \tilde{t} = W(c_0, \hat{c}_0) t.$$

Remark 5.20. Dividing by the Wronskian, which is constant, is predicated on the absence of 0 in the spectrum of the original problem. For Neumann-Neumann boundary conditions 0 is an eigenvalue and $W(c_0, \hat{c}_0) = 0$. Then one possible choice is $c_0 = \hat{c}_0 = \frac{1}{\sqrt{2}}$ and $b_{-1} = \frac{1}{2}, b_0(x) = \int_0^1 G_T(x, \xi) \rho(\xi) d\xi$, where $G_T(x, y) = -\frac{|x-y|}{2}$ is the translationally invariant Green's function of $-D_x^2$.

We will finish this section stating explicit formulas for the evolution of positions x_i and weights m_i in the case of the discrete measure ρ . These formulas are implicit in the statement of the previous theorem and we present them only to emphasize their simplicity.

Corollary 5.21. Let the ZS deformation be given by $b = b_0 + \frac{b_{-1}}{z}$, with $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, $0 < x_1 < \dots < x_N < 1$ and let b_0 and b_{-1} be given by equations (5.24). Then the positions x_j and masses m_j evolve with respect to the rescaled time $\tilde{t} = W(c_0, \hat{c}_0) t$ according to:

$$\dot{x}_i = -b_0(x_i), \quad \dot{m}_i = m_i \langle b_{0,x} \rangle(x_i), \quad 1 \leq i \leq N,$$

and the sums

$$I_1 = \sum_{i=1}^N m_i G(x_i, x_i),$$

and

$$I_n = \sum_{1 \leq i_1 < i_2 < \dots < i_n \leq N} m_{i_1} m_{i_2} \dots m_{i_n} (x_{i_n} - x_{i_{n-1}})(x_{i_{n-1}} - x_{i_{n-2}}) \dots (x_{i_2} - x_{i_1}) G(x_{i_1}, x_{i_n}), \quad 2 \leq n \leq N$$

are invariants of the ZS deformation (constants of motion).

6 Examples of isospectral deformations of a discrete string

Example 6.1. As an example of the developed formalism we consider the case of the string boundary value problem (1.2) corresponding to the mass density $\rho = \sum_{j=1}^N m_j \delta_{x_j}$, $0 < x_1 < x_2 < \dots < x_n < 1$, with boundary conditions of the type $0 < h < \infty$ and $H = 0$. Let

$$v|_{I_j} = v_j = p_j(x - x_j) + q_j, \quad \text{where } I_j = (x_j, x_{j+1}),$$

denote the solution to the initial value problem $-v_{xx} = z\rho v$, $v_x(0) - hv(0) = 0$ with the convention that $x_0 = 0$ and $x_{N+1} = 1$. The construction of this solution proceed as follows.

From the boundary condition at $x = 0$, one can take $p_0 = h$ and $q_0 = 1$. Letting l_j denote the length of the interval I_j and imposing continuity of v at $x = x_{j+1}$ one finds that $p_j = \frac{q_{j+1} - q_j}{l_j}$. The

jump in the derivative of v at $x = x_{j+1}$ results in $p_{j+1} - p_j = -zm_{j+1}q_{j+1}$. On the last interval I_N $v_N = A_1(x-1) + B_1 = p_N(x-x_N) + q_N$ from which one finds that $A_1 = p_N$ and $B_1 = p_N l_N + q_N$. Let us define the Weyl function for this problem

$$W(z) \stackrel{\text{def}}{=} \frac{v(1; z)}{v_x(1; z)} = \frac{B_1}{A_1}.$$

From the construction $W(z) = \frac{p_N l_N + q_N}{p_N} = l_N + \frac{q_N}{p_N}$. Iterating with the help of continuity and jump conditions we obtain

$$W(z) = l_N + \frac{1}{-zm_N + \frac{1}{l_{N-1} + \frac{1}{\ddots + \frac{1}{l_0 + \frac{1}{h}}}}}.$$

Since the spectrum of the boundary value problem at hand is simple (only simple eigenvalues occur) we have the following partial fraction decomposition: $W(z) = W_\infty + \sum_{i=1}^N \frac{a_i}{-z+z_i}$, where $W_\infty = l_N$ and

$$a_i = - \lim_{z \rightarrow z_i} \frac{B_1(z)}{A_1(z)} (z - z_i) = - \frac{B_1(z_i)}{\frac{dA_1}{dz} |_{z=z_i}}.$$

It can be shown that the residues a_i are strictly positive. We recall that the spectrum $\{z_j : 1 \leq j \leq N\}$ is determined from $A_1(z) = 0$. Following Theorem 5.16, item (8), we get $\dot{A}_1 = 0$, and $\dot{B}_1(z_i, t) = 2\beta(z_i)B_1(z_i, t)$ so that

$$B_1(z_i, t) = B_1(z_i, t=0) e^{2\beta(z_i)t}.$$

This implies that the residues evolve as

$$a_i(t) = a_i(0) e^{2\beta(z_i)t}.$$

Thus the Weyl function can be written as

$$W(z, t) = l_N + \sum_{i=1}^N \frac{a_i(0) e^{2\beta(z_i)t}}{-z+z_i} \stackrel{\text{def}}{=} l_N + \int \frac{e^{2\beta(\zeta)t} d\mu(\zeta)}{-z+\zeta}$$

where $d\mu(\zeta) = \sum_{i=1}^N a_i(0) \delta_{z_i}$ is the spectral measure at $t=0$. Hence if we set $\lambda = -z$ we obtain the formula

$$\int \frac{e^{2\beta(\zeta)t} d\mu(\zeta)}{\lambda+\zeta} = \frac{1}{\lambda m_N + \frac{1}{l_{N-1} + \frac{1}{\ddots + \frac{1}{l_0 + \frac{1}{h}}}}},$$

which shows that using Stieltjes' inversion formulas (see [5, 24]) we can recover m_N, \dots, m_1 and l_{N-1}, \dots, l_0 in terms of the moments of the measure $e^{2\beta(\zeta)t} d\mu(\zeta)$. Finally, one can recover l_N by observing that

$$W(0) = \frac{h+1}{h} = 1 + \frac{1}{h} = l_N + \int \frac{e^{2\beta(\zeta)t} d\mu(\zeta)}{\zeta}$$

and solve for l_N , or compute l_N from the formula for the total length of the string: $l_N = 1 - \sum_{i=0}^{N-1} l_i$. The formula for $\beta(\zeta)$ depends on the deformation and boundary conditions and is given by (5.22), (5.24), if $\epsilon > 0$, $\epsilon = 0$ respectively. More generally, for the multi-pole model, $\beta(\zeta)$ is given by (5.14).

Example 6.2. Let us now consider the case of $\epsilon = 0$, hence $b = b_0 + \frac{b_{-1}}{z}$, and let us assume a discrete mass density $\rho = \sum_{j=1}^N m_j \delta_{x_j}$. For the sake of comparison we consider two cases of boundary conditions: Dirichlet-Dirichlet and Dirichlet-Neumann. We point out that the Dirichlet-Neumann case appeared, somewhat unexpectedly, in the recent work on the modified Camassa-Holm equation [10].

We use the formulas corresponding to the rescaled time (see theorem 5.19 and corollary 5.21).

For the Dirichlet-Neumann case we have

$$G^{DN}(x, y) = - \begin{cases} x, & x < y \\ y, & x > y \end{cases}$$

$$b_{-1} = -x, \quad b_0 = - \int_0^1 |x - \xi| G^{DN}(x, \xi) \rho(\xi) d\xi = \sum_{x_j < x} (x - x_j) x_j m_j + \sum_{x_j > x} (x_j - x) x m_j,$$

Both b_{-1}, b_0 satisfy the required boundary conditions $b(0) = 0, b_{xx}(1) = 0$.

For the Dirichlet-Dirichlet case we have

$$G^{DD}(x, y) = - \begin{cases} x(1 - y), & x < y \\ y(1 - x), & x > y \end{cases}$$

$$b_{-1} = -x(1 - x), \quad b_0 = - \int_0^1 |x - \xi| G^{DD}(x, \xi) \rho(\xi) d\xi = \sum_{x_j < x} (x - x_j)(1 - x) x_j m_j + \sum_{x_j > x} (x_j - x)x(1 - x_j) m_j,$$

Clearly, b_{-1}, b_0 satisfy the required boundary conditions $b(0) = 0, b(1) = 0$.

In both cases, the explicit form of the evolution equations (with respect to the rescaled time) reads:

$$\dot{x}_i = \sum_{j=1}^N |x_i - x_j| G(x_i, x_j) m_j,$$

$$\dot{m}_i = -m_i \left(\sum_{j=1}^N \operatorname{sgn}(x_i - x_j) G(x_i, x_j) m_j + \sum_{j=1}^N |x_i - x_j| \langle G_x(x, x_j) \rangle (x_i) m_j \right), \quad \operatorname{sgn}(0) = 0.$$

with G being an appropriate Green's function.

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8 Appendix I: weak Lax pairs

We assume $\rho = \sum_{j=1}^N m_j \delta_{x_j}$ for the remainder of the appendix. Our goal is to compute conditions on the Lax pair to be compatible as a distributional Lax pair. We will denote by D_x, D_t the distributional derivatives with respect to x, t respectively, and define the distributional Lax pair as:

$$D_x V = \begin{bmatrix} 0 & 1 \\ -z\rho & 0 \end{bmatrix} V \quad D_t V = \begin{bmatrix} -\frac{1}{2}D_x(b) + \beta & b \\ -\frac{1}{2}D_x^2(b) - z\rho b & \frac{1}{2}D_x(b) + \beta \end{bmatrix} V. \quad (8.1)$$

It is immediate from the first equation that the first component of V is continuous and piecewise smooth with jumps only on the support of ρ . So we can write, as before, $V = \begin{bmatrix} v \\ v_x \end{bmatrix}$, where v_x is the classical derivative. This implies that the first member of the Lax pair contains only one non-trivial statement: $D_x(v_x) = D_x^2(v) = -z\rho v$ which is equivalent to

$$v_{xx} = 0, \quad \text{if } x \neq x_j, \quad [v_x](x_j) = -zm_j v(x_j), \quad (8.2)$$

where $[f](x_j)$ means the jump of f at x_j . Likewise, the second member of the Lax pair is well defined if b is continuous and piecewise smooth with jumps only on the support of ρ . Thus the second equation can be written:

$$D_t V = \begin{bmatrix} v_t \\ v_{tx} - \sum_{i=1}^N [v_x](x_i) \dot{x}_i \delta_{x_i} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2}D_x(b) + \beta & b \\ -\frac{1}{2}D_x^2(b) - z\rho b & \frac{1}{2}D_x(b) + \beta \end{bmatrix} V, \quad (8.3)$$

or, equivalently,

$$v_t = \left(-\frac{1}{2}b_x + \beta\right)v + bv_x, \\ v_{tx} - \sum_{i=1}^N [v_x](x_i) \dot{x}_i \delta_{x_i} = -\frac{1}{2}b_{xx}v - \sum_{j=1}^N \left(\frac{1}{2}[b_x](x_j)v(x_j) + zm_j b(x_j)v(x_j)\right) \delta_{x_j} + \left(\frac{1}{2}b_x + \beta\right)v_x,$$

showing that it is the singular part of $D_t V$ that determines the evolution of x_i ; using (8.2) and (8.3) we obtain

$$zm_i \dot{x}_i = -\left(\frac{1}{2}[b_x](x_i) + zm_i b(x_i)\right). \quad (8.4)$$

The actual compatibility condition of equations (8.1), (8.3), which can be succinctly formulated as distributional equality $D_x D_t V = D_t D_x V$, is presented in the following theorem.

Theorem 8.1. *Let $\rho = \sum_{j=1}^N m_j \delta_{x_j}$ and let b be continuous, piecewise differentiable function with jumps only on the support of ρ . Then $D_x D_t V = D_t D_x V$ implies:*

$$zD_t \rho = \frac{1}{2}D_x^3 b + z\mathcal{L}_\rho b, \quad (8.5)$$

where $\mathcal{L}_\rho b = D_x(\rho b) + \langle b_x \rangle \rho$ and $\langle b_x \rangle$ means the pointwise, arithmetic average of the right hand and left hand limits.

Proof. First we observe that the left hand side contains only the singular part of the distributional equality. Indeed,

$$D_t \rho = \sum_{i=1}^N \dot{m}_i \delta_{x_i} - \sum_{i=1}^N m_i \dot{x}_i \delta'_{x_i}.$$

Hence the regular part of the right hand side must be zero. Equating coefficients at δ_{x_j} and δ'_{x_j} on both sides we obtain that the statement of the theorem is equivalent to:

$$b_{xxx} = 0, \quad \text{if } x \neq x_i, \quad (8.6a)$$

$$z \dot{m}_i = \frac{1}{2} [b_{xx}](x_i) + z m_i \langle b_x \rangle(x_i) \quad (8.6b)$$

$$-z m_i \dot{x}_i = \frac{1}{2} [b_x](x_i) + z m_i b(x_i), \quad (8.6c)$$

for $i = 1, \dots, N$. Moreover by localizing tests functions around individual points x_i we can, without loss of generality, assume that $N = 1$; in other words we can work with $\rho = m_1 \delta_{x_1}$.

We obtain two statements corresponding to components of $D_t D_x V = D_x D_t V$:

$$v_{tx} - \dot{x}_1 [v_x](x_1) \delta_{x_j} = \left(-\frac{1}{2} b_x + \beta \right) v + b v_x \Big|_x + \left[-\frac{1}{2} b_x v + b v_x \right](x_1) \delta_{x_1} \quad (8.7a)$$

$$\begin{aligned} & -z(\dot{m}_1 v(x_1) + m_1 \dot{v}(x_1)) \delta_{x_1} + z m_1 v(x_1) \dot{x}_1 \delta'_{x_1} = \quad (8.7b) \\ & \left(-\frac{1}{2} b_{xx} v + \frac{1}{2} b_x v_x + \beta v_x \right) \Big|_x + \left[-\frac{1}{2} b_{xx} v + \frac{1}{2} b_x v_x + \beta v_x \right](x_1) \delta_{x_1} - \left(\frac{1}{2} [b_x](x_1) v(x_1) + z m_1 b(x_1) v(x_1) \right) \delta'_{x_1}. \end{aligned}$$

The equality of regular parts implies:

$$v_{tx} = \left(-\frac{1}{2} b_x + \beta \right) v + b v_x \Big|_x = v_{xt}, \quad b_{xxx} = 0,$$

while equating in the first equation the coefficients at δ_{x_1} , with the help of (8.2), one easily recovers (8.6c). Finally, to prove (8.6b) we need to equate the coefficients of δ_{x_1} :

$$-z(\dot{m}_1 v(x_1) + m_1 \dot{v}(x_1)) = \left[-\frac{1}{2} b_{xx} v + \frac{1}{2} b_x v_x + \beta v_x \right](x_1),$$

and proceed by observing that

1. $\dot{v}(x_1) = \langle v_x \rangle \dot{x}_1 + \langle v_t \rangle(x_1)$,
2. $[b_x v_x] = \langle b_x \rangle [v_x] + [b_x] \langle v_x \rangle$.

Upon carrying out several elementary cancellations one arrives, somewhat exhausted, at the final formula (8.6b). \square

9 Appendix II: mapping to the Real Line

In this appendix we discuss a relation between the flow generated by $b = b_0 + \frac{b_{-1}}{z}$ and the CH equation [9] (see also [4]), which we will write in the notation adapted to our setup:

$$m_t = \mathcal{L}_m u, \quad m = \frac{1}{2}(u_{\zeta\zeta} - u), \quad \zeta \in \mathbf{R}, \quad (9.1)$$

where $\mathcal{L}_m = m \frac{d}{d\zeta} + \frac{d}{d\zeta} m$. We will discuss the relation for the smooth case only, although the comparison is also valid on the level of distributional Lax pairs. We recall the Lax pair formulation of CH:

$$-\Psi_{\zeta\zeta} + \frac{1}{4}\Psi = zm\Psi, \quad \Psi_t = (u - \frac{1}{z})\Psi_{\zeta} - \frac{u_{\zeta}}{2}\Psi, \quad (9.2)$$

whose compatibility condition gives the evolution equation on m given above and the constraint

$$m_{\zeta} = \frac{1}{2}(u_{\zeta\zeta} - u)_{\zeta}. \quad (9.3)$$

For us the starting equations are (2.11)

$$\rho_t = \mathcal{L}_{\rho} b_0, \quad \frac{1}{2}b_{0,xxx} + \mathcal{L}_{\rho} b_{-1} = 0, \quad b_{-1,xxx} = 0,$$

where ρ is the mass density in the boundary value problem undergoing an isospectral deformation

$$-v_{xx} = z\rho v, \quad v_t = (-\frac{1}{2}b_x + \beta)v + bv_x. \quad (9.4)$$

Now we perform a Liouville transformation:

$$x = \frac{1}{2} + \frac{1}{2} \tanh \frac{\zeta}{2}, \quad \Psi = (\cosh \frac{\zeta}{2}) v, \quad (9.5)$$

which leads to an equivalent form of Lax equations (9.4) on the real axis:

$$-\Psi_{\zeta\zeta} + \frac{1}{4}\Psi = z(\frac{1}{16}(\operatorname{sech}^4 \frac{\zeta}{2})\rho)\Psi, \quad \Psi_t = 4(\cosh^2 \frac{\zeta}{2})b\Psi_{\zeta} + (-\frac{1}{2}((4\cosh^2 \frac{\zeta}{2})b)_{\zeta} + \beta)\Psi, \quad (9.6)$$

If we define:

$$\hat{u} = 4(\cosh^2 \frac{\zeta}{2})b \stackrel{\text{def}}{=} \hat{u}_0 + \frac{\hat{u}_{-1}}{z}, \quad \hat{m} = \frac{1}{16}(\operatorname{sech}^4 \frac{\zeta}{2})\rho \quad (9.7)$$

then the compatibility condition for (9.6) reads:

$$\hat{m}_t = \mathcal{L}_{\hat{m}} \hat{u}_0, \quad (9.8a)$$

$$\frac{1}{2}(\hat{u}_0 - \hat{u}_{0,\zeta\zeta})_{\zeta} = \mathcal{L}_{\hat{m}} \hat{u}_{-1}, \quad (9.8b)$$

$$(\hat{u}_{-1} - \hat{u}_{-1,\zeta\zeta})_{\zeta} = 0. \quad (9.8c)$$

For the Dirichlet-Dirichlet case (see Example 6.2)

$$b_{-1} = -x(1-x) = -\frac{1}{4(\cosh^2 \frac{\zeta}{2})}$$

which implies $\hat{u}_{-1} = -1$, which by equation (9.8b) implies the constraint (9.3), and one recovers the CH case. On the other hand, for example, for the Dirichlet-Neumann case

$$b_{-1} = -x$$

and hence

$$\hat{u}_{-1} = -(e^\zeta + 1)$$

which does not reduce to the CH case.

References

- [1] M. P. Appell. Sur la transformation des equations differentielles lineaires. *Comptes Rendus des Seances*, 91:211–214, 1880.
- [2] V. Barilon. Inverse eigenvalue problems. In *Inverse problems (Montecatini Terme, 1986)*, volume 1225 of *Lecture Notes in Math.*, pages 1–51. Springer, Berlin, 1986.
- [3] J. H. Barrett. Canonical forms for third-order linear differential equations. *Ann. Mat. Pura Appl.* (4), 65:253–274, 1964.
- [4] R. Beals, D. H. Sattinger, and J. Szmigielski. Acoustic scattering and the extended Korteweg-de Vries hierarchy. *Adv. Math.*, 140(2):190–206, 1998.
- [5] R. Beals, D. H. Sattinger, and J. Szmigielski. Multi-peakons and a theorem of Stieltjes. *Inverse Problems*, 15(1):L1–L4, 1999.
- [6] R. Beals, D. H. Sattinger, and J. Szmigielski. Inverse scattering solutions of the Hunter-Saxton equation. *Appl. Anal.*, 78(3-4):255–269, 2001.
- [7] R. Beals, D. H. Sattinger, and J. Szmigielski. The string density problem and the Camassa-Holm equation. *Philos. Trans. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci.*, 365(1858):2299–2312, 2007.
- [8] G. Borg. Eine Umkehrung der Sturm-Liouvilleschen Eigenwertaufgabe. Bestimmung der Differentialgleichung durch die Eigenwerte. *Acta Math.*, 78:1–96, 1946.
- [9] R. Camassa and D. D. Holm. An integrable shallow water equation with peaked solitons. *Phys. Rev. Lett.*, 71(11):1661–1664, 1993.
- [10] X. Chang, X. Hu, and J. Szmigielski. Multipeakons of a two-component modified Camassa-Holm equation and the relation with the finite Kac-van Moerbeke lattice. *Submitted*, 2015.
- [11] H. Dym and H. P. McKean. *Gaussian processes, function theory, and the inverse spectral problem*. Academic Press [Harcourt Brace Jovanovich Publishers], New York, 1976. Probability and Mathematical Statistics, Vol. 31.
- [12] L. D. Faddeev and L. A. Takhtajan. *Hamiltonian methods in the theory of solitons*. Springer Series in Soviet Mathematics. Springer-Verlag, Berlin, 1987. Translated from the Russian by A. G. Reyman [A. G. Reĭman].

- [13] G. M. L. Gladwell. *Inverse problems in vibration*, volume 119 of *Solid Mechanics and its Applications*. Kluwer Academic Publishers, Dordrecht, second edition, 2004.
- [14] H. P. W. Gottlieb. Isospectral strings. *Inverse Problems*, 18(4):971–978, 2002.
- [15] T. Helmuth, R. Spiteri, and J. Szmigielski. One-dimensional magnetotelluric inversion with radiation boundary conditions. *Can. Appl. Math. Q.*, 15(4):419–445, 2007.
- [16] J. K. Hunter and R. Saxton. Dynamics of director fields. *SIAM J. Appl. Math.*, 51(6):1498–1521, 1991.
- [17] E. L. Ince. *Ordinary Differential Equations*. Dover Publications, New York, 1944.
- [18] I. S. Kac and M. G. Krein. On the spectral functions of the string. *Amer. Math. Soc. Transl.*, 103(2):19–102, 1974.
- [19] M. G. Krein. On inverse problems for a nonhomogeneous cord. *Doklady Akad. Nauk SSSR (N.S.)*, 82:669–672, 1952.
- [20] J. Pöschel and E. Trubowitz. *Inverse spectral theory*, volume 130 of *Pure and Applied Mathematics*. Academic Press Inc., Boston, MA, 1987.
- [21] P. C. Sabatier. On some spectral problems and isospectral evolutions connected with the classical string problem. I. Constants of motion. *Lett. Nuovo Cimento (2)*, 26(15):477–482, 1979.
- [22] P. C. Sabatier. On some spectral problems and isospectral evolutions connected with the classical string problem. II. Evolution equation. *Lett. Nuovo Cimento (2)*, 26(15):483–486, 1979.
- [23] P. C. Sabatier. Around the classical string problem. In *Nonlinear evolution equations and dynamical systems (Proc. Meeting, Univ. Lecce, Lecce, 1979)*, volume 120 of *Lecture Notes in Phys.*, pages 85–102. Springer, Berlin, 1980.
- [24] T.-J. Stieltjes. Recherches sur les fractions continues. *Ann. Fac. Sci. Toulouse Sci. Math. Sci. Phys.*, 8(4):J1–J122, 1894.
- [25] S. Tefteller. Concerning solutions of third order self-adjoint differential equations. *Ann. Mat. Pura Appl. (4)*, 96:185–192, 1972.
- [26] P. Weidelt. The Inverse Problem of Geomagnetic Induction. *Z. f. Geophysik*, 38:257–289, 1970.